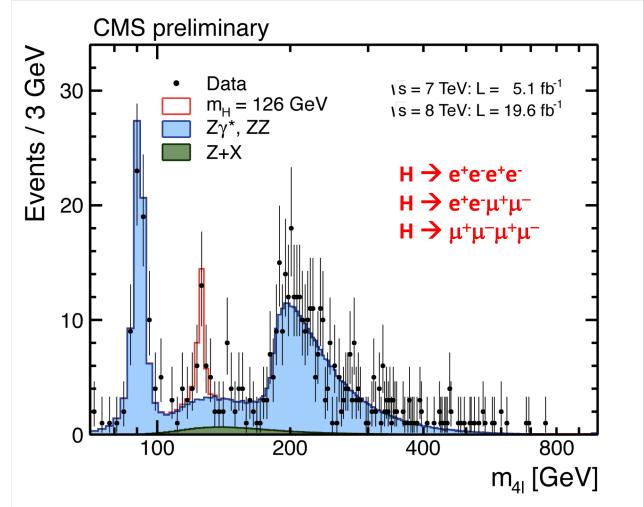
A common problem in physics

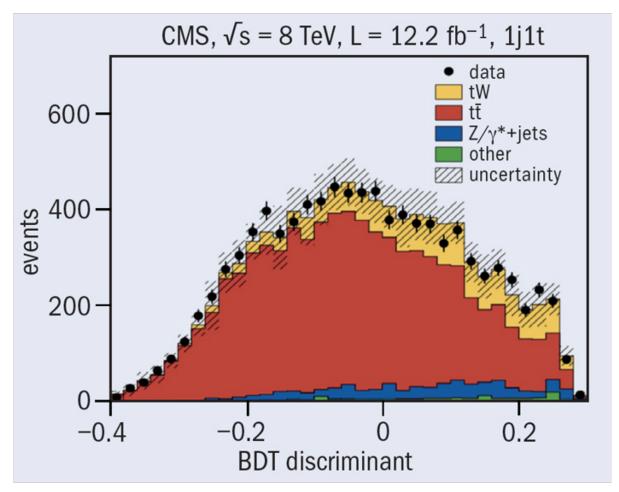
- Dig out "signal" for some interesting but rare process out of a sea of "background"
- Toy example here from high energy physics
- First step: select "events" with patterns of final state particles/kinematics according to what your signal should look like to distinguish it from background
- Most likely this selection will give you both signal and background events
- You then have to
 - Decide whether you have a signal or not
 - If so, estimate how "strong" your signal is
 - because you are trying to "measure" something physical, eg, a cross-section

Sometimes it is fairly obvious



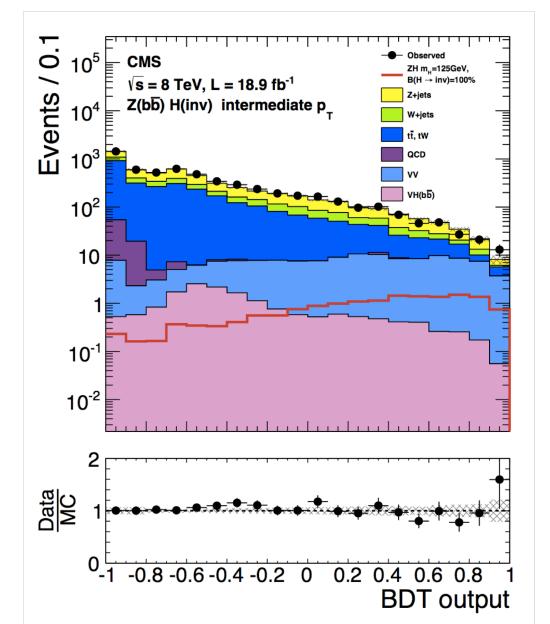
Obviously there is a signal. Precisely estimating the strength of the signal is a different story

Sometimes less so



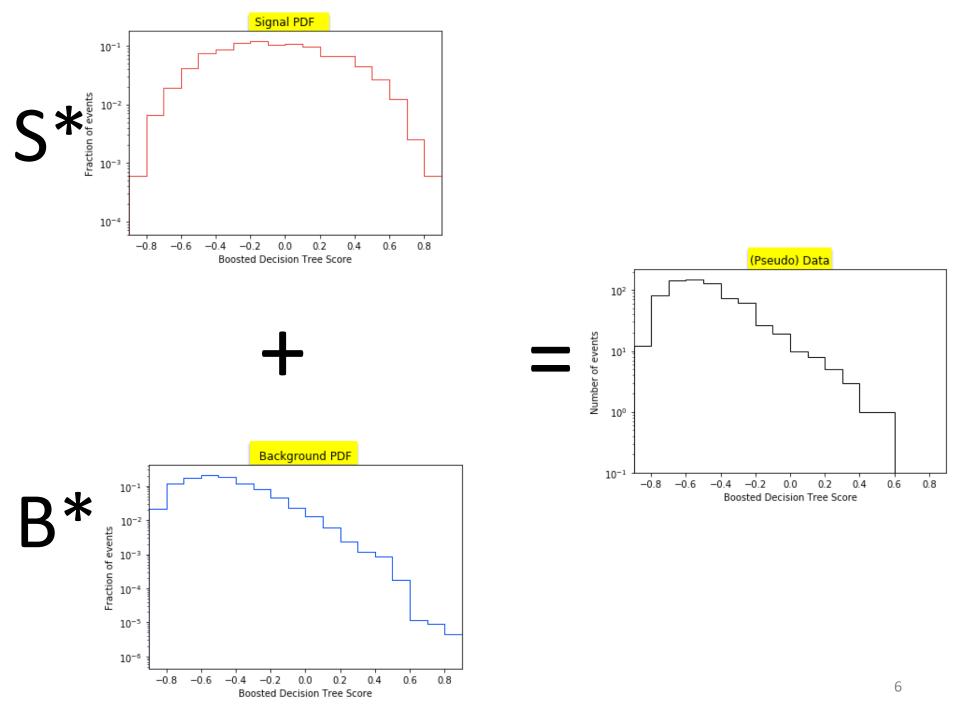
Here the signal is " $pp \rightarrow tW$ ". (t=top quark...W=W boson) What is plotted is a "machine learning" "boosted decision tree" discriminant which combines a lot of information into one single variable that looks different for signal and background

Sometimes it's kind of hopeless



- In general have to fit the data distribution to sum of signal (S) + background (B) distributions
- This could mean fitting
 - A single 1D histogram
 - A 2D (or 3D, or..) histogram
 - Several different histograms simultaneously
 - Or even single events in an <u>unbinned</u> fashion
- Obviously you need to select and treat your data intelligently, and understand the pdfs for S and B

 Including uncertainties
- We now look at a semi-realistic example



- Want to estimate number of signal (S) and background events (B)
 - <u>actually, we usually mostly only care about S</u>
 - See "Appendix" for what we really mean by this
- Use the histograms as PDFs
- Because many of the bins are low stats, need
 Poisson Uncertainties → use NLL
- $\{s_i\}$ and $\{b_i\}$ are the <u>binned</u> pdfs for S and B

- Normalized to 1, ie, $\Sigma s_i = 1$ and $\Sigma b_i = 1$

- $\{d_i\}$ are contents of i-th bin in the (pseudo) data
- The model for the {d_i} is

$$\mu_i = S^*s_i + B^*b_i$$

$$\mathcal{L} = \prod p(d_i) = \prod \frac{e^{-\mu_i} \mu_i^{d_i}}{d_i!}$$

$$\begin{aligned} &-\log \mathcal{L} &= \sum \mu_i - d_i \log \mu_i \\ &-\log \mathcal{L} &= \sum S \cdot s_i + B \cdot b_i - d_i \log(S \cdot s_i + B \cdot b_i) \\ &-\log \mathcal{L} &= S + B - \sum d_i \log(S \cdot s_i + B \cdot b_i) \end{aligned}$$

This is actually called "extended log likelihood" (I dropped an additive constant from the NLL) Note: it does not look at all like a χ^2 . But as was discusses previously in the large d_i limit it is the same as χ^2 within a factor of $\frac{1}{2}$

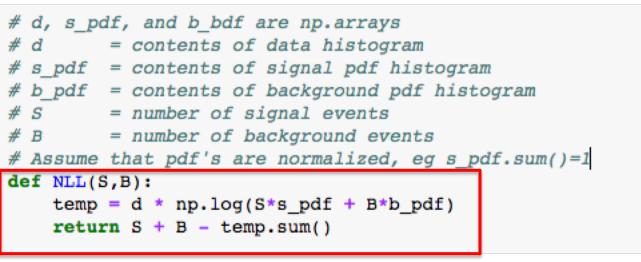
We will fit with Minuit

iminuit tutorial at

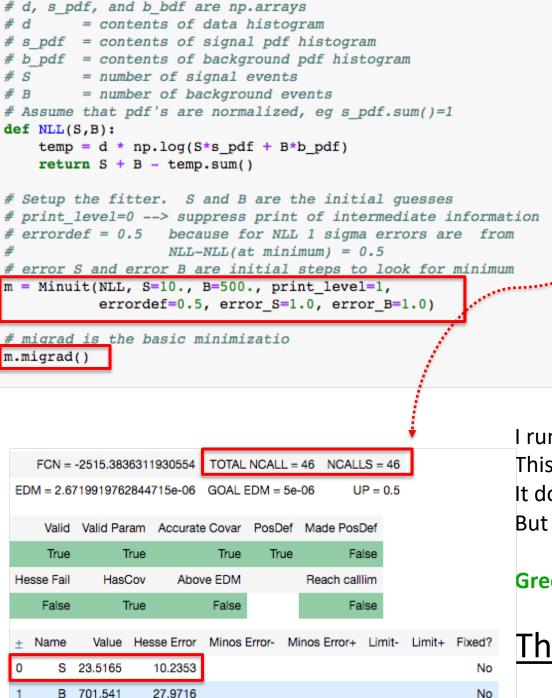
https://nbviewer.jupyter.org/github/iminuit/iminuit/blob/master/tutorial/basic_tutorial.ipynb

import numpy as np import matplotlib.pyplot as plt from iminuit import Minuit

Define the negative log likelihood function



/home/pi/physrpi/campagnari/python/maxLikFit.py



Note: I could have told Minuit to impose the constraint S>0 and/or B>0

This is best to be avoided for technical reasons.

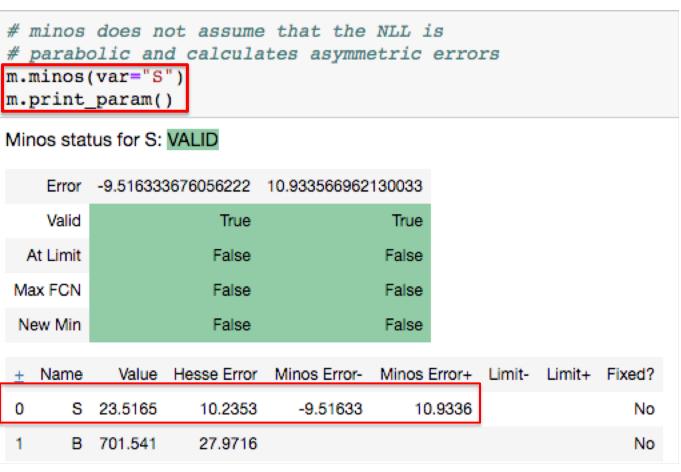
However: with S and/or B < 0 some of the calls to NLL could result in taking log of negative number. I should probably have protected against it, but luckily it did not happen ⁽²⁾

I run the code through a "jupyter notebook" This gives a nicely formatted output. It does not look as nice when run normally. But you get the same info.

Green means "it worked"

<u>The answer is <u>S = 23.5 ± 10.2</u></u>

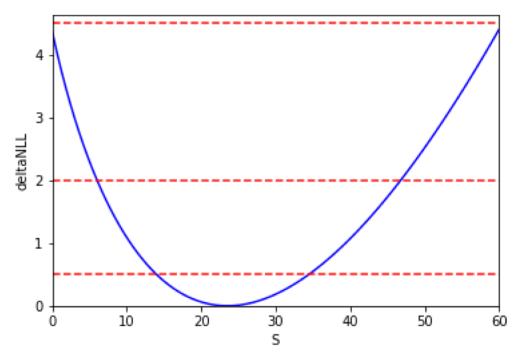
More sophisticated analysis: "minos"



 $S = 23.5^{+10.9}_{-9.5}$

PS: the dataset had 700 background events and 25 signal events₁₁

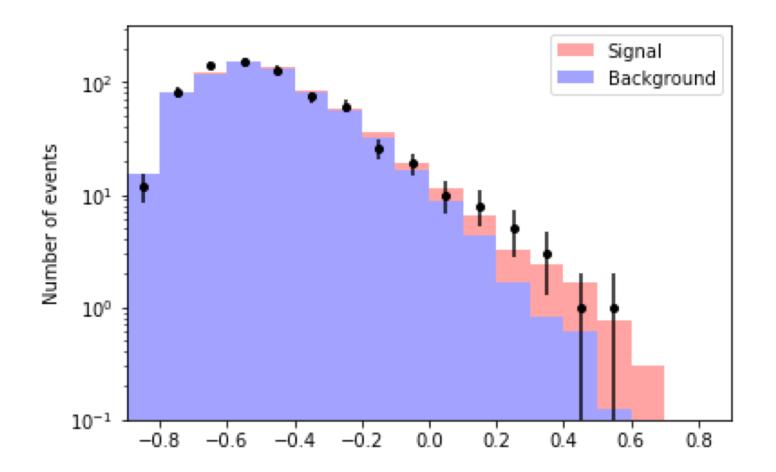




Can see by eye that it is not exactly parabolic.

The asymmetry around the minimum is what gives asym. minos errors

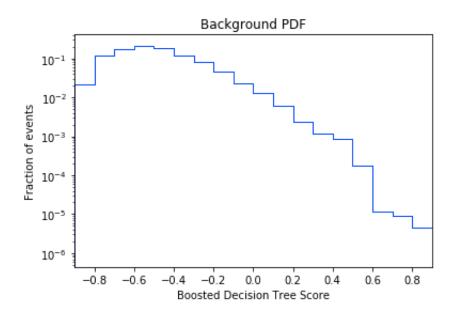
This is almost 3σ away from 0, but not quite.

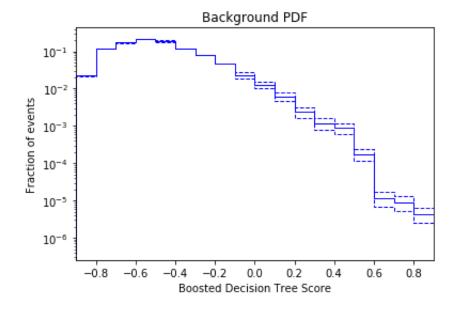


Data here is plotted with sqrt(N) uncertainties. Not quite right, but conventional

Systematics (simple example)

Imagine we do not know the background pdf perfectly





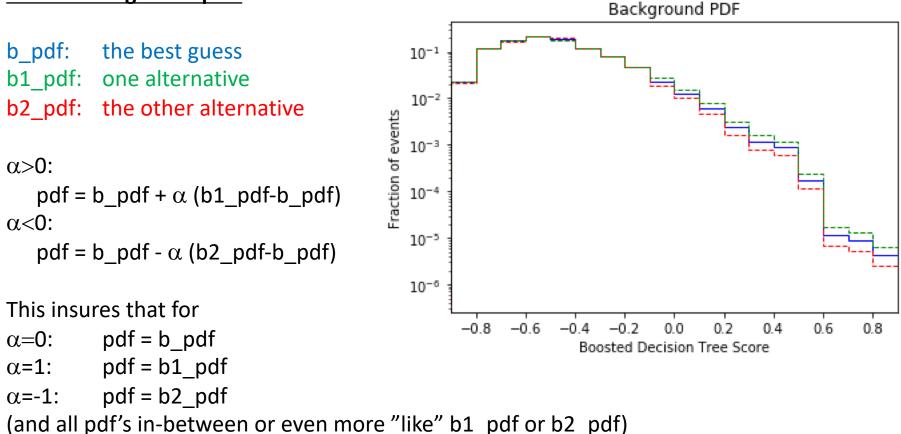
This is our best guess for the bg pdf (what we have used so far).

Comes from some separate study we may have done

The dashed lines "bracket" at the 1 sigma level our lack of knowledge of the bg pdf. This "information" also would come from some study

How to take this "shape uncertainty" into account?

Introduce a new parameter α that smoothly interpolates the options for the background pdf:



Our best guess for α is $\alpha=0$, with $\sigma=1$. Let the data itself figure out what α should be Modify the likelihood to include information on α α becomes a 3rd parameter to be fit for Include "prior" information on α in likelihood "Nuisance parameter"

$$\mathcal{L}
ightarrow \mathcal{L} \ e^{-rac{lpha^2}{2}} \ -\log \mathcal{L}
ightarrow -\log \mathcal{L} + rac{lpha^2}{2^{15}}$$

```
# d, s pdf, and b bdf are np.arrays
                                                      # d, s pdf, b bdf, b1 pdf, b2 pdf are np.arrays
                                                      # d = contents of data histogram
# d = contents of data histogram
                                                      # s pdf = contents of signal pdf histogram
# s pdf = contents of signal pdf histogram
                                                      # b pdf = contents of default background pdf histogram
# b pdf = contents of background pdf histogram
                                                      # b1 pdf = contents of alternative 1 to b pdf
# S = number of signal events
                                                      # b2 pdf = contents of alternative 2 too b pdf
# B = number of background events
                                                      # S = number of signal events
# Assume that pdf's are normalized, eg s pdf.sum()=1
                                                      # B = number of background events
def NLL(S,B):
                                                      # alpha = parameter to interpolate between pdfs
   temp = d * np.log(S*s pdf + B*b pdf)
                                                      # Assume that pdf's are normalized, eg s pdf.sum()=1
   return S + B - temp.sum()
                                                      def new NLL(S,B,alpha):
                                                          # code below insures that
                                                          # alpha=0 ---> use b pdf
                                                          # alpha=1 ---> use b1 pdf
                                                          # alpha=-1 ---> use b2 pdf
                                                          # (and smoothly interpolates vs. alpha)
                                                          if alpha>0:
                                                             new_b_pdf = b_pdf + alpha*(b1_pdf-b_pdf)
                                                          else:
                                                             new b pdf = b pdf - alpha*(b2 pdf-b pdf)
                                                          # should be already normalized, but make sure
                                                          new b pdf = new b pdf / new b pdf.sum()
                                                          temp = d * np.log(S*s_pdf + B*new_b_pdf)
                                                          return S + B - temp.sum() + alpha*alpha/2.
```

```
# d, s_pdf, and b_bdf are np.arrays
# d = contents of data histogram
# s_pdf = contents of signal pdf histogram
# b_pdf = contents of background pdf histogram
# S = number of signal events
# B = number of background events
# Assume that pdf's are normalized, eg s_pdf.sum()=1
def NLL(S,B):
   temp = d * np.log(S*s_pdf + B*b_pdf)
   return S + B - temp.sum()
```

Shows the flexibility of Minuit. The function to minimize is a NLL with a bunch of Poisson terms from data counts, but also with an additional term that has nothing to do with data counts and is most certainly not Poissonian

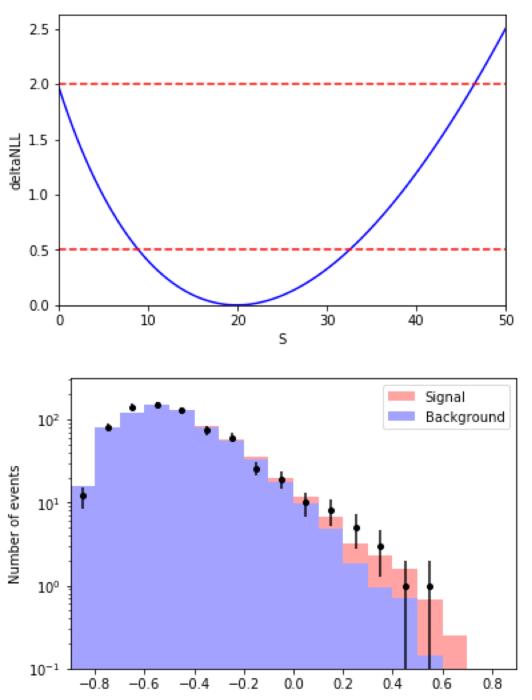
```
# d, s pdf, b bdf, b1 pdf, b2 pdf are np.arrays
        = contents of data histogram
# d
# s pdf = contents of signal pdf histogram
# b pdf = contents of default background pdf histogram
# b1 pdf = contents of alternative 1 to b pdf
# b2 pdf = contents of alternative 2 too b pdf
# S
      = number of signal events
# B
      = number of background events
# alpha = parameter to interpolate between pdfs
# Assume that pdf's are normalized, eg s pdf.sum()=1
def new NLL(S,B,alpha):
    # code below insures that
    # alpha=0 ---> use b pdf
    # alpha=1 ---> use b1 pdf
    # alpha=-1 ---> use b2 pdf
    # (and smoothly interpolates vs. alpha)
    if alpha>0:
       new_b_pdf = b_pdf + alpha*(b1_pdf-b_pdf)
    else:
       new_b_pdf = b_pdf - alpha*(b2_pdf-b_pdf)
    # should be already normalized, but make sure
    new b pdf = new b pdf / new b pdf.sum()
    temp = d * np.log(S*s_pdf + B*new_b_pdf)
    return S + B - temp.sum() + alpha*alpha/2.
```

Min	ios stat	us for S: V	ALID	Re	esult	ts w	vith	shap	be ur
	Error	-11.025370	084171879	12.6085760132	21337				
	Valid		True		True				
At Limit			False						
Max FCN			False		False				
New Min			False		False				
±	Name	Value	Hesse Error	Minos Error-	Minos	Error+	Limit-	Limit+	Fixed?
0	S	19.9292	11.8446	-11.0254	12	2.6086			No
1	В	705.074	28.7197						No
2	alpha	-0.412914	0.746692						No
			<i>S</i> =	19.9	$+1 \\ -1$	2.6 1 0			

used to be

 $S = 23.5^{+10.9}_{-9.5}$

Uncertainty increased because we told the fitter that we are not entirely sure about the shape of the background (makes sense that we lose accuracy)



Results with shape unc.

Just about 2σ from zero

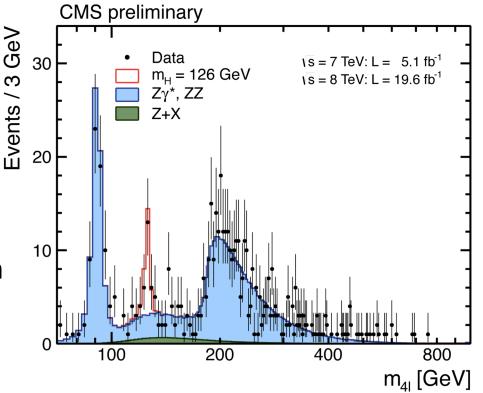
If the signal (or the background!) varies very fast, it makes sense to make the bins small enough to capture the important features of the distribution <u>CMS preliminary</u>

In this example the signal (in red) is sharply peaked. The binsize is *small enough*.

A binsize of 50 GeV would have completely lost all the information on the signal.

Could make binsize even smaller

In the limit that we let the binsize go to zero, we can perform an <u>unbinned</u> ML fit



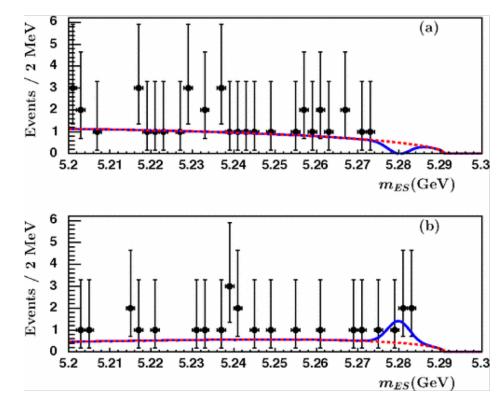
Example of unbinned fit results

The S+B fit is blue The B component is red

This is an interesting case because the fitted S in the top panel is negative!

 $S = -1.6^{+0.7}_{-0.0}$ Note: the fit is unbinned but the

data are plotted binned (how else could we plot it?)



In order to perform an unbinned fit it is best if the the pdf's for S and B are continuous functions (instead of histograms) But how does the procedure that we outlined change? Binned fit:

$$-\log \mathcal{L} = S + B - \sum d_i \log(S \cdot s_i + B \cdot b_i)$$

 d_i = number of data counts in the *i*-th bin

 $S s_i$ = number of expected counts in the *i*-th bin from signal

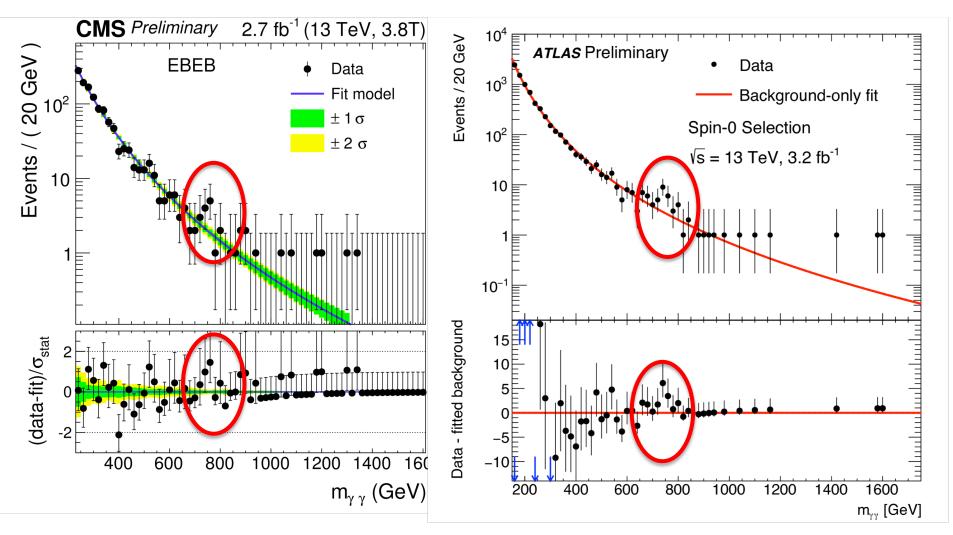
B b_i = number of expected counts in the *i*-th bin from background The sum is over <u>bins</u>

Unbinned fit:

$$-\log \mathcal{L} = S + B - \sum \log(S \cdot s(x_i) + B \cdot b(x_i))$$

 x_i = measured value of discriminating variable for *i-th* events(x)= pdf for signalb(x)= pdf for backgroundThe sum is over eventsYou will doing an
unbinned fit in a
future homework

A surprise: diphoton mass bump near 750 GeV?



Data from 1st higher energy LHC run (May-Nov 2015) presented in mid-December 2015

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How the $\gamma\gamma$ Resonance Stole Christmas

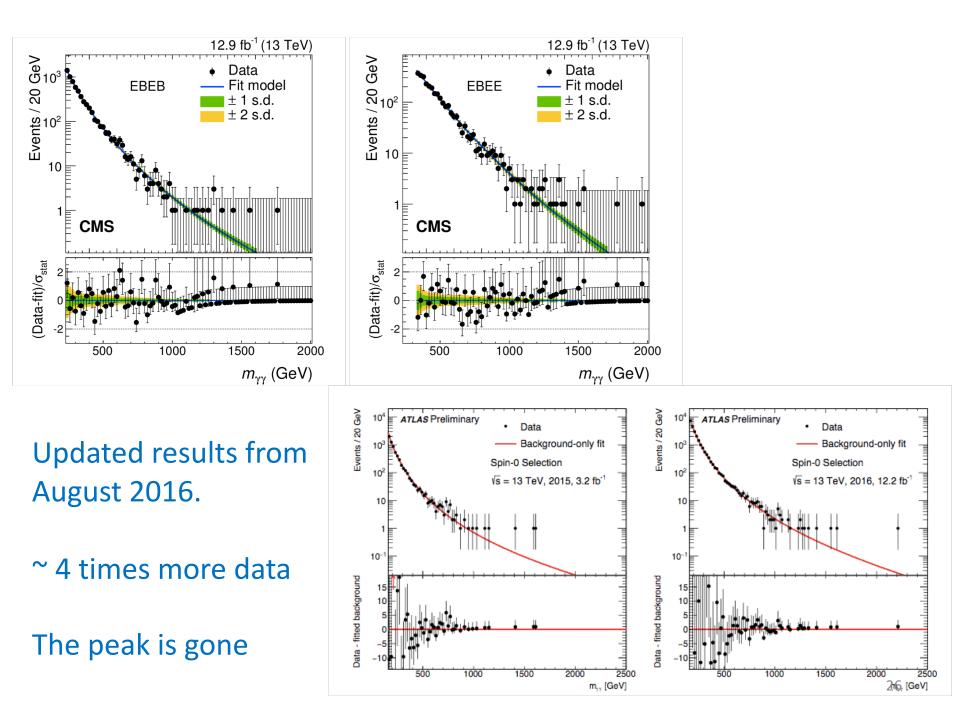
Nathaniel Craig[♠], Patrick Draper[♠], Can Kilic[♦], and Scott Thomas[♣]

Department of Physics, University of California, Santa Barbara, CA 93106, USA

- Weinberg Theory Group, Department of Physics and Texas Cosmology Center, The University of Texas at Austin, Austin, TX 78712, USA
- * New High Energy Theory Center, Rutgers University, Piscataway, NJ 08854, USA

Abstract

The experimental and theoretical implications of heavy di-gauge boson resonances that couple



APPENDIX

Technicality

- The uncertainty on S returned by the extended log likelihood fit is <u>no</u>t the uncertainty on the number of signal events contained in the particular dataset
- This fact can be understood with a simple example
 - Suppose B=0.
 - Only 1 bin
 - Let N = the number of events
 - Since B=0, S=N
 - There is no uncertainty whatsoever <u>on the number of</u> <u>signal events contained in this dataset</u>. I have N events, there is no background, so all of them are signal. Full stop.
- Let's see what the formalism actually gives us, see next page....

$$-\log \mathcal{L}(S, B) = S + B - \sum d_i \log(S \cdot s_i + B \cdot b_i)$$

$$-\log \mathcal{L}(S, B = 0) = S - d_1 \log(S \cdot s_1)$$

$$-\log \mathcal{L}(S) = S - N \log S$$

This is minimized for *S*=*N* as expected.

The change in NLL moving away from the minimum S=N by δN is

 $\Delta \text{NLL} = -\log \mathcal{L}(N + \delta N) + \log \mathcal{L}(N)$ $\Delta \text{NLL} = \delta N - N \log(1 + \frac{\delta N}{N})$

Expanding the log for small $\delta N/N$ and large N):

$$\Delta \text{NLL} \approx \delta N - N(\frac{\delta N}{N} - \frac{1}{2}(\frac{\delta N}{N})^2) = \frac{\delta N^2}{2N}$$

Thus, the "1 sigma" uncertainty δN that one obtains by setting $\Delta NLL = \frac{1}{2}$ is: $\delta N = \sigma = \sqrt{N}$

This is the usual counting uncertainty in the Gaussian regime that you can interpret in a frequentist or baysean sense, as you like.