

# Compact Notation

The Levi-Civita symbol in three dimension,  $\epsilon_{ijk}$ , where  $i$ ,  $j$ , and  $k$  are integers between 1 and 3, is defined as:

- $\epsilon_{ijk} = 0$  if any of  $i$ ,  $j$ , and  $k$  are the same. For example  $\epsilon_{122} = 0$ .
- $\epsilon_{ijk} = +1$  if  $ijk$  is an **even** permutation of 123. So  $\epsilon_{123} = \epsilon_{312} = \epsilon_{231} = +1$ .
- $\epsilon_{ijk} = -1$  if  $ijk$  is an **odd** permutation of 123. So  $\epsilon_{132} = \epsilon_{321} = \epsilon_{123} = -1$ .

Another convention that leads to a more compact notation is that **sum over repeated indices is implied**.

For example, imagine that we have vectors  $\vec{a}$  and  $\vec{b}$  in 3D with components  $(a_x, a_y, a_z) = (a_1, a_2, a_3)$  and  $(b_x, b_y, b_z) = (b_1, b_2, b_3)$ . Then we could write

$$\begin{aligned} |\vec{a}|^2 &= \sum_i a_i a_i = a_i a_i \\ \vec{a} \cdot \vec{b} &= \sum_i a_i b_i = a_i b_i \\ (\vec{a} \times \vec{b})_k &= \sum_{i,j} \epsilon_{ijk} a_i b_j = \epsilon_{ijk} a_i b_j. \end{aligned}$$

The Levi-Civita symbol allows you to write cross-products in a very compact way. If you do not believe me, write out explicitly the components of  $\vec{a} \times \vec{b}$  and check for yourself.

Furthermore, the commutation relations of the angular momentum operator can be written in compact form as

$$[L_i, L_j] = i\epsilon_{ijk} \hbar L_k$$

As you can see, the  $\epsilon_{ijk}$  factor insures that the commutator of any component of angular momentum with itself ( $i = j$ ) is zero, as it should; further,  $[L_i, L_j] = -[L_j, L_i]$  since  $\epsilon_{ijk} = -\epsilon_{jik}$ ; finally, the only non-zero term in the implied sum over  $k$  on the right side is the one with  $k \neq i$  and  $k \neq j$ , meaning that the commutator of two different components of  $\vec{L}$  gives  $\pm\hbar$  times the third component, with the sign determined by the order of the two elements of the commutator.