

DIRAC NOTATION, bras & kets

We have used $\psi(x)$ ie wavefunctions to represent state of QM system

We have also written $|\psi\rangle$

More abstract \equiv

Also used for spin states $|\uparrow\rangle$

There is also correspondence $\psi(x)$

$$\psi(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \psi(x) e^{-ipx/\hbar} dx$$

position space \longleftrightarrow momentum space

"change of representation"

Finite vector space

$$\vec{a} = a_i \hat{e}_i = \sum a_i \hat{e}_i = \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix}$$

$$\vec{a}' = a'_i \hat{e}'_i = \begin{pmatrix} a'_1 \\ \vdots \\ a'_n \end{pmatrix}$$

$|a\rangle$ Abstract way of representing state
"ket"

eigenvectors of $\hat{x} |x\rangle = x |x\rangle$

eigenvectors of $\hat{p} |p\rangle = p |p\rangle$

eigenvector of $\hat{S}_z |\uparrow\rangle = +\frac{1}{2}\hbar |\uparrow\rangle$

~~✓~~

Inner product of two vectors

$|a\rangle, |b\rangle$

$\langle b|a\rangle$

$\langle b|a\rangle$ What is $\langle b|$ "bra"

$|a\rangle$ is "ket" represents a state

Bra is not a vector

One to one correspondence $|b\rangle \leftrightarrow \langle b|$

Mathematically speaking:

— ket lives in a vector space V

— bras are elements of a dual space of linear functionals on V

linear functionals

$\langle b|$ acts on $|a\rangle \rightarrow$ number

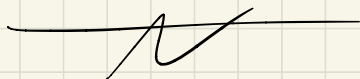
$\langle b|a\rangle =$ complex numbers

Operator on ket \rightarrow another ket

$$\hat{O} |a\rangle = |c\rangle$$

Functional on ket \rightarrow number

$$\langle b|a\rangle = \text{number}$$



Finite V

$$|a\rangle = \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} \quad |b\rangle = \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix}$$

$$\langle b|a\rangle$$

$$\langle b| = (b_1^* \dots b_n^*)$$

$$\langle b|a\rangle = (b_1^* \dots b_n^*) \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} =$$

$$\langle b|a\rangle = b_i^* a_i$$

$$\langle b| = (|b\rangle)^\dagger$$

OUTER PRODUCT

$$|a\rangle\langle b| \quad ???$$

finite V

$$|a\rangle\langle b| = \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} (b_1^* \dots b_n) = \begin{pmatrix} N \times N \\ \text{matrix} \end{pmatrix}$$

$|a\rangle\langle b| \Rightarrow$ is an operator

$$\underline{\underline{(|a\rangle\langle b|)}} \underline{\underline{|c\rangle}} = |a\rangle \underbrace{\langle b|c\rangle}_{\text{number}} = \underline{\underline{\text{number } |a\rangle}}$$

Let's denote basis $|\phi_i\rangle$ orthonormal

$$|\psi\rangle = c_i |\phi_i\rangle$$

$$\langle \phi_j | \psi \rangle = c_i \langle \phi_j | \phi_i \rangle = c_i \delta_{ij} = c_j$$

$$\rightarrow |\psi\rangle = \underbrace{\langle \phi_i | \psi \rangle}_{= c_i} |\phi_i\rangle$$

$$|\psi\rangle = \sum_i |\phi_i\rangle \underbrace{\langle \phi_i | \psi \rangle}_{\text{operator}}$$

$$|\psi\rangle = \left(\sum_i |\phi_i\rangle \langle \phi_i| \right) |\psi\rangle$$

= Identity

$$\sum_i |\phi_i\rangle \langle \phi_i| = \mathbb{1}$$

completeness

For a continuous set $|\phi(i)\rangle$

$$\int |\phi(i)\rangle \langle \phi(i)| di = \underline{1}$$

i is whatever labels the states,
eg p or x or ...

$|\psi\rangle$ to $\psi(x)$

$$\psi(x) \equiv \langle x | \psi \rangle$$

$$|\psi\rangle = \int_{-\infty}^{\infty} dx' |x'\rangle \langle x' | \psi \rangle =$$

$$\langle x | \psi \rangle = \int_{-\infty}^{\infty} dx' \underbrace{\langle x | x' \rangle}_{\delta(x-x')} \langle x' | \psi \rangle$$

$$\underline{\underline{\langle x | \psi \rangle}} = \int_{-\infty}^{\infty} dx' \underbrace{\langle x | x' \rangle}_{\delta(x-x')} \underbrace{\langle x' | \psi \rangle}_{\text{number}}$$

$$\langle x | \psi \rangle = \psi(x)$$

$$\psi(x) = \int_{-\infty}^{\infty} dx' \delta(x-x') \psi(x')$$

\bar{x} = expectation value of x

$$\bar{x} = \int_{-\infty}^{\infty} x |\psi(x)|^2 dx$$

$$\bar{x} = \langle \psi | \hat{x} | \psi \rangle$$

$$\bar{x} = \int dx' \langle \psi | \underbrace{\hat{x} | x' \rangle}_{x' | x' \rangle} \langle x' | \psi \rangle$$

$$\bar{x} = \int dx' x' \langle \psi | x' \rangle \langle x' | \psi \rangle = \int dx' x' |\langle x' | \psi \rangle|^2$$

$$\psi(x) = \langle x | \psi \rangle$$

Symmetries, Conservation laws

Spatial Translation (1D)

Classically: move particle by a

QM: take state and shift it by a
in space

$T(a)$ = operator which shifts
state by a

$$T(a)|n\rangle = |n+a\rangle$$

$$T(a)|\psi\rangle = |\phi\rangle$$

$$\psi(x) = \langle n|\psi\rangle$$

$$\phi(x) = \langle n|\phi\rangle = \langle n|T(a)|\psi\rangle$$

$$\phi(x) = \langle u | T(a) | \psi \rangle$$

$$\phi(x) = \int dx' \langle u | T(a) | x' \rangle \langle x' | \psi \rangle$$

$\underbrace{\hspace{10em}}_{|x'+a\rangle} \quad \underbrace{\hspace{10em}}_{\psi(x')}$

$$\phi(x) = \int dx' \langle x | x'+a \rangle \psi(x')$$

$\underbrace{\hspace{10em}}_{\delta(x-x'-a)}$

$\hookrightarrow \underbrace{x' = x - a}$

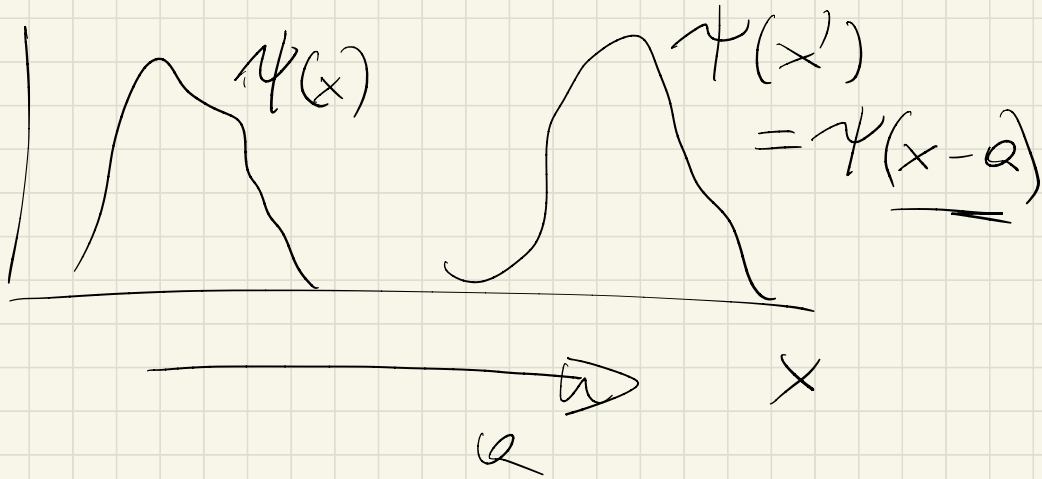
$$\phi(x) = \psi(x-a)$$

$$T(a) |u\rangle = |u+a\rangle$$

$$\psi(x) \rightarrow \psi(x-a)$$

$$|x\rangle \rightarrow |x+a\rangle$$

$$\psi(x) \rightarrow \psi(x-a)$$



$$T(0) = \underline{1}$$

$$T(a)T(b) = T(a+b)$$

$$T(a)^{-1} = T(-a)$$

$$T(a)^{-1} = T(a)^{\dagger} \quad \text{unitary}$$

$$\langle \phi | \phi \rangle = 1 \quad \langle \psi | \psi \rangle = 1$$

$$\phi = T(a) \psi$$

$$\langle T(a) \psi | T(a) \psi \rangle = 1$$

$$\langle \psi | \underbrace{T(a)^\dagger T(a)}_{=1} | \psi \rangle = 1$$

$T(a)$ not hermitian

$$T(a) = e^{-\frac{i a \hat{P}}{\hbar}}$$