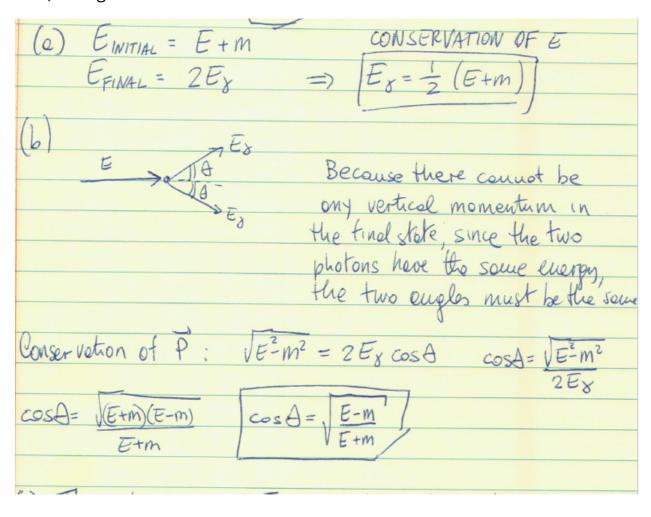
Physics 110B Practice Final Solutions

Problem 1

Working with c=1. In the end result, to restore the right units, if you care about that, change m into mc^2 .



(a) The induced current I a must be such as to counteract
maining the change in flux - Since the solenoid
field is to the right and is decreasing I
field is to the right and is decreasing Ic must be such that it makes a B field to the
pointing to the right
Using the right-hond-rule, this means Ic
pointing to the right- Using the right-head-rule, this means Ic gos from LEFT TO RIGHT
(b) 0=5BT(c)=5 TdB
$1 \qquad 2/4 \qquad D(t) = 577 \text{ MeT} (t)$
(b) $\phi = 5BT(\frac{d}{2})^2 = \frac{5}{4}\pi d^3B$ $\beta = 4TI(N/2) \implies \phi(t) = \frac{5}{1}\pi d^3I(t)$ $E = \frac{5}{1}\pi d^3B$
5 turns C

Induced emf

$$I_{c} = \underbrace{E}_{R} = \underbrace{I}_{C} \underbrace{-1}_{D} \underbrace{d\Phi}_{R} = \underbrace{I}_{R} \underbrace{-1}_{C} \underbrace{-1}_{R} \underbrace{I}_{R} \underbrace{-1}_{R} \underbrace{-1}_{R} \underbrace{I}_{R} \underbrace{-1}_{R} \underbrace{-1}_{R}$$

Displacement current
$$J = \mathcal{E}_0 \frac{\partial \mathcal{E}}{\partial t}$$
 (FCR)

 $J_D = \frac{\mathcal{E}_0 \alpha}{\pi R^2 \sigma}$ Independent of r for $r \in \mathbb{R}$. Independent of t (b) Maxwell eqtn $\sqrt{x} \hat{B} = M_0 \hat{J} + M_0 \mathcal{E}_0 \frac{\partial \hat{\mathcal{E}}}{\partial t}$
 $\sqrt{x} \hat{B} = M_0 (\hat{J} + \hat{J}_0)$

When turning this into an integral equation this becomes $\int \hat{B} d\hat{\ell} = M_0 \hat{I}_{enclosed} + \hat{I}_{Denclosed}$

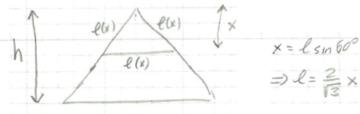
Take path as circle of radius $r > R$

Also $I_{Denclosed} = J_{D} \cdot \pi R^2 = \mathcal{E}_0 \alpha$
 $B = \frac{M_0 \alpha}{2\pi r} \left\{ t + \frac{\mathcal{E}_0}{\sigma} \right\}$

(a)
$$\mathcal{E} = -\frac{d\Phi}{dt}$$
 $I = \frac{\mathcal{E}}{R}$

All angles in the triangle are 60°-=> Height of triengle h = a sin 60° = 13 a

Picture of triengle



=) l= = x

B(x) = Mo I (x+b is the distance from the wire) Flux through a swell horizontal strip et x: do= B(x) ldx But since $\ell = \frac{2}{\sqrt{3}} \times i$ $d\phi = 2B \times j \times dx = hoI \times dx$ Total Plux $\phi = \frac{\mu_0 I}{\sqrt{2} \pi} \int_{-x+b}^{x} dx = \frac{\mu_0 I}{\sqrt{2} \pi} \int_{-x+b}^{x} dx$ Φ = MOI Sh dx - MOID Sh dx P= MOIR - MOID ROG h+b And since h= \(\frac{13}{2}a\) \phi = \(\frac{\pmoI}{1}\) \(\frac{a}{2}\) - \(\frac{b}{13}\) \(\lambda_2\) \(\frac{1}{2}a\) + \(\frac{b}{13}\) And finelly M= of = 40 [a - b log = +b]

The speed of the wave is:

$$v = \frac{1}{\sqrt{\epsilon \mu_o}} = \frac{1}{\sqrt{f(x)\mu_o}}$$

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But also

$$v = \frac{\mathrm{d}x}{\mathrm{d}t} \Rightarrow \mathrm{d}x = v \, \mathrm{d}t \Rightarrow \mathrm{d}t = \frac{\mathrm{d}x}{v}$$

With the speed from the previous relationship,

$$dt = \sqrt{\mu_o} \sqrt{f(x)} dx$$

The time needed for the wave to pass the layer of thickness *l* is found by integration:

$$t = \sqrt{\mu_o} \int_0^l \sqrt{f(x)} \, \mathrm{d}x$$

Now let us obtain the constants A and b for function f(x) from the boundary conditions.

$$f(0) = \epsilon_1, f(l) = \epsilon_2, \epsilon_2 < \epsilon_1$$

 $f(0) = A = \epsilon_1, f(l) = Ae^{-bl} = \epsilon_2$

From the second equation it follows that $\epsilon_1 e^{-bl} = \epsilon_2$,

$$e^{-bl} = \frac{\epsilon_2}{\epsilon_1} \Rightarrow -bl = \ln \frac{\epsilon_2}{\epsilon_1} \Rightarrow b = -\frac{1}{l} \ln \frac{\epsilon_2}{\epsilon_1}$$

The time is, therefore,

$$t = \sqrt{\mu_o} \int_0^l \sqrt{Ae^{-bx}} \, dx = \sqrt{\mu_o A} \left(-\frac{2}{b} e^{-\frac{bx}{2}} \right) \Big|_0^l = \frac{2}{b} \sqrt{\mu_o A} \left(1 - e^{-\frac{bl}{2}} \right)$$

Back to the constants $A = \epsilon_1$, $b = -\frac{1}{l} \ln \frac{\epsilon_2}{\epsilon_1}$ and

$$t = \frac{2}{-\frac{1}{l} \ln \frac{\epsilon_2}{\epsilon_1}} \sqrt{\mu_o \epsilon_1} \left(1 - e^{-\frac{1}{2} \ln \frac{\epsilon_2}{\epsilon_1}} \right) = \frac{2l \sqrt{\mu_o}}{\ln \frac{\epsilon_2}{\epsilon_1}} (\sqrt{\epsilon_2} - \sqrt{\epsilon_1}) = \frac{2l \sqrt{\mu_o}}{\ln \frac{\epsilon_1}{\epsilon_2}} (\sqrt{\epsilon_1} - \sqrt{\epsilon_2})$$

The motion of the electron involves energy that can be broken into three pieces: energy associated with moving a charged particle though a potential difference, kinetic energy, and energy to account for radiative losses

$$U_{\text{pot}} = qV$$
 $U_{\text{kin}} = \frac{1}{2}mv^2$ $U_{\text{rad}} = P_{\text{rad}}t$

We will consider $v \ll c$ so the Larmor formula can be used.

$$P_{\rm rad} = \frac{\mu_o q^2 a^2}{6\pi c}$$

The initial kinetic energy is zero and the final is given by

$$U_{\rm kin} = \frac{1}{2} m_e v^2 = \frac{1}{2} m_e (2ad) = m_e ad$$

where we use

$$v^2 = v_o^2 + 2ad$$

with $v_0 = 0$. For the potential energy, q = -e and $V = -V_0$, so

$$U_{\text{pot}} = (-e)(-V_o) = eV_o$$

For the radiation energy, we have

$$P_{\rm rad} = \frac{\mu_o q^2 a^2}{6\pi c} = \frac{\mu_o e^2 a^2}{6\pi c}$$

and considering,

$$a = \frac{v}{t} \Rightarrow t = \frac{v}{a} = \frac{\sqrt{2ad}}{a} = \sqrt{\frac{2d}{a}}$$

we have

$$U_{\rm rad} = P_{\rm rad}t = \frac{\mu_o e^2 a^2}{6\pi c} \sqrt{\frac{2d}{a}} = \frac{\mu_o e^2 a^{3/2}}{3\pi c} \sqrt{\frac{d}{2}}$$

The total energy is then given by

$$\begin{split} U &= U_{\rm kin} + U_{\rm pot} + U_{\rm rad} = m_e a d + e V_o + \frac{\mu_o e^2 a^{3/2}}{3\pi c} \sqrt{\frac{d}{2}} \\ &= a d \left(m_e + \frac{\mu_o e^2}{3\pi c} \sqrt{\frac{a}{2d}} \right) + e V_o \end{split}$$