SESSION 3

GRIFFITHS 8.28 (a) lach q has $\vec{v} = wR\hat{\phi} + v_{\vec{z}}\hat{z}$ $\vec{F}_{i} = q\left(\vec{v} \times \vec{B}\right) = qk\left(2wRz\hat{p} - Rv_{z}\hat{\phi} + wR^{2}\hat{z}\right)$ Mon each charge _ Net force is $F = \sum_{i=1}^{n} F_i$ As we go around the circle, the p and of forces concel out, and all we are left with is the 2 components $nqRWR^{2} = M \frac{d^{2}}{d^{2}} \hat{z}$ So we have $\frac{1}{2} - \frac{nqkR^2}{m}w(t)$

In order to solve this we need w(f) Consider now the torque $\overline{N} = \sum_{i=1}^{n} \overline{\Gamma_{i} \times F_{i}} = n R \hat{\rho} \times (-g k R \sigma_{z}) \hat{\phi}$ $\overline{N} = -nqkR^2 \overline{D_2} = I \frac{dw}{dt} \hat{z} \qquad I = moment$ This then gives dw _ ngkR² d2 dt I dt $\frac{d^2}{d\omega} = -\frac{nqkR^2}{dz} \frac{d^2}{dz}$ df2 I Then using the equation circled in block: $\frac{d^2 w}{dt^2} = - \left(\frac{nqkr^2}{NI}\right)^2 w$ oud since J= MR

 $\frac{d^2 \omega}{dt^2} = -2 \frac{n^2 q^2 k^2 \hat{R}}{M^2} \omega$ Let $\beta^2 = 2n^2 q^2 k_r R^2 = W(t) = W_p \cos \beta t$ Going bock to the equation circled in red $\frac{dz}{dt} = -\frac{T}{nqkR^2} \frac{dw}{dt} = -\frac{t'_2 mR^2}{nqkR^2} \beta w_0 \left(-\sin\beta t\right)$ $\frac{dz}{dt} = \frac{1}{\sqrt{2}} \frac{R}{M} w_0 \sin\beta t$ $\frac{dz}{dt} = \frac{1}{\sqrt{2}} \frac{R}{M} w_0 \sin\beta t$ $Z(t) = \frac{R_{Wo}}{\Gamma_2 M} \left(1 - \cos \beta t \right)$ 4 Z(0)=0 Check conservation of every $\frac{1}{2}M\mathcal{T}_{z}(t) + \frac{1}{2}\mathcal{I}\tilde{w}(t) = E$ But $N_2(t) = \frac{dz}{dt} = \frac{1}{\sqrt{z}} \frac{R}{M}$ wo singt

 $E = \frac{1}{2} \frac{M}{2M^2} \frac{R^2}{W_0^2} \frac{W_0^2 \sin^2 \beta \epsilon}{\sin^2 \beta \epsilon} + \frac{1}{2} \frac{I}{2} \frac{W_0^2 \cos^2 \beta \epsilon}{\cos^2 \beta \epsilon}$ $E = \frac{1}{2} \frac{MR^2}{2} w_0 \sin^2 \beta t + \frac{1}{2} \left(\frac{MR^2}{2}\right) w_0^2 \cos^2 \beta t$ $E = \frac{1}{2} \frac{MR^2}{2} \omega_0^2 = \frac{1}{2} T \omega_0^2$