

Srednicki Chapter 38

QFT Problems & Solutions

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Srednicki 38.1. Use equation 38.12 to compute $u_s(\mathbf{p})$ and $v_s(\mathbf{p})$ explicitly. Hint: Show that the matrix $2i\hat{\mathbf{p}} \cdot \mathbf{K}$ has eigenvalues ± 1 and that, for any matrix A with eigenvalues ± 1 , $e^{cA} = \cosh c + (\sinh c)A$, where c is an arbitrary complex number.

We have:

$$2iK^j = -\gamma^i \gamma^0$$

$$2iK^j = - \begin{pmatrix} 0 & \gamma^j \\ -\gamma^j & 0 \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} = \begin{pmatrix} -\sigma^j & 0 \\ 0 & \sigma^j \end{pmatrix}$$

At the risk of stating the obvious, note that this represents three matrices, each one indicated by j . Each of these matrices has eigenvalues ± 1 . Any linear combination (of unit magnitude) of these matrices will similarly have eigenvalues ± 1 . Thus, $2i\hat{\mathbf{p}} \cdot \mathbf{K}$ will be a single matrix with eigenvalues ± 1 .

Next, we simply expand:

$$\exp(cA) = \sum_{n \text{ even}} \frac{(cA)^n}{n!} + \sum_{n \text{ odd}} \frac{(cA)^n}{n!}$$

Working in a basis where A is diagonal, we have $A^2 = 1$ (recall that we defined A to have eigenvalues ± 1). Then:

$$\exp(cA) = \sum_{n \text{ even}} \frac{(c)^n}{n!} + \sum_{n \text{ odd}} \frac{(c)^n}{n!} A$$

This gives:

$$\exp(cA) = \cosh c + (\sinh c)A$$

Combining these results, we have:

$$\exp(i\eta\hat{\mathbf{p}} \cdot \mathbf{K}) = \cosh\left(\frac{\eta}{2}\right) + \sinh\left(\frac{\eta}{2}\right) \begin{pmatrix} -\hat{\mathbf{p}} \cdot \vec{\sigma} & 0 \\ 0 & \hat{\mathbf{p}} \cdot \vec{\sigma} \end{pmatrix}$$

We can write this as one matrix:

$$\exp(i\eta\hat{\mathbf{p}} \cdot \mathbf{K}) = \begin{pmatrix} \cosh(\eta/2) - \sinh(\eta/2)\hat{\mathbf{p}} \cdot \vec{\sigma} & 0 \\ 0 & \cosh(\eta/2) + \sinh(\eta/2)\hat{\mathbf{p}} \cdot \vec{\sigma} \end{pmatrix} \quad (38.1.1)$$

To simplify this, we need to do something with $\hat{p} \cdot \vec{\sigma}$. Let's expand:

$$\begin{aligned}
\hat{p} \cdot \vec{\sigma} &= p_x \sigma_1 + p_y \sigma_2 + p_z \sigma_3 \\
&= p \sin \theta \cos \phi \sigma_1 + p \sin \theta \sin \phi \sigma_2 + p \cos \theta \sigma_3 \\
&= p \sin \theta \cos \phi \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + p \sin \theta \sin \phi \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + p \cos \theta \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\
&= p \sin \theta \begin{pmatrix} 0 & \cos \phi - i \sin \phi \\ \cos \phi + i \sin \phi & 0 \end{pmatrix} + p \cos \theta \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\
&= p \sin \theta \begin{pmatrix} 0 & e^{-i\phi} \\ e^{i\phi} & 0 \end{pmatrix} + p \cos \theta \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\
&= p \begin{pmatrix} \cos \theta & \sin \theta e^{-i\phi} \\ \sin \theta e^{i\phi} & -\cos \theta \end{pmatrix}
\end{aligned}$$

Combining this with equation (38.1.1), we have:

$$\exp(i\eta\hat{p}\cdot K) = \begin{pmatrix} \cosh(\eta/2) - p \sinh(\eta/2) \cos \theta & -p \sinh(\eta/2) \sin \theta e^{-i\phi} & 0 & 0 \\ -p \sinh(\eta/2) \sin \theta e^{i\phi} & \cosh(\eta/2) + p \sinh(\eta/2) \cos \theta & 0 & 0 \\ 0 & 0 & \cosh(\eta/2) + p \sinh(\eta/2) \cos \theta & p \sinh(\eta/2) \sin \theta e^{-i\phi} \\ 0 & 0 & p \sinh(\eta/2) \sin \theta e^{i\phi} & \cosh(\eta/2) - p \sinh(\eta/2) \cos \theta \end{pmatrix}$$

Now we just have to multiply some matrices, according to equation 38.12. The result is:

$$u_+(\vec{p}) = \sqrt{m} \begin{pmatrix} \cosh(\eta/2) - p \sinh(\eta/2) \cos \theta \\ -p \sinh(\eta/2) \sin \theta e^{i\phi} \\ \cosh(\eta/2) + p \sinh(\eta/2) \cosh \theta \\ p \sinh(\eta/2) \sin \theta e^{i\phi} \end{pmatrix}$$

$$u_-(\vec{p}) = \sqrt{m} \begin{pmatrix} -p \sinh(\eta/2) \sin \theta e^{-i\phi} \\ \cosh(\eta/2) + p \sinh(\eta/2) \cos \theta \\ p \sinh(\eta/2) \sin \theta e^{-i\phi} \\ \cosh(\eta/2) - p \sinh(\eta/2) \cos \theta \end{pmatrix}$$

$$v_+(\vec{p}) = \sqrt{m} \begin{pmatrix} -p \sinh(\eta/2) \sin \theta e^{-i\phi} \\ \cosh(\eta/2) + p \sinh(\eta/2) \cos \theta \\ -p \sinh(\eta/2) \sin \theta e^{-i\phi} \\ -\cosh(\eta/2) + p \sinh(\eta/2) \cos \theta \end{pmatrix}$$

$$v_-(\vec{p}) = \sqrt{m} \begin{pmatrix} -\cosh(\eta/2) + p \sinh(\eta/2) \cos \theta \\ p \sinh(\eta/2) \sin \theta e^{i\phi} \\ \cosh(\eta/2) + p \sinh(\eta/2) \cosh \theta \\ p \sinh(\eta/2) \sin \theta e^{i\phi} \end{pmatrix}$$

Srednicki 38.2. Verify equation 38.15.

We have:

$$\gamma^0 = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$\gamma^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}$$

So:

$$\gamma^{0\dagger} = \gamma^0$$

$$\gamma^{i\dagger} = -\gamma^i$$

since the Pauli matrices are Hermitian.

Then we have:

$$\overline{\gamma^0} = \beta\gamma^0\beta$$

$$\overline{\gamma^i} = -\beta\gamma^i\beta$$

which is:

$$\overline{\gamma^0} = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$\overline{\gamma^i} = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} 0 & -\sigma^i \\ \sigma^i & 0 \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

Doing the multiplication, we have:

$$\overline{\gamma^0} = \gamma^0$$

$$\overline{\gamma^i} = \gamma^i$$

as expected. □

Next for $S^{\mu\nu}$:

$$\overline{S^{\mu\nu}} = \beta(S^{\mu\nu})^\dagger\beta$$

We have:

$$S^{\mu\nu} = \frac{i}{4} \begin{pmatrix} \sigma^\mu\bar{\sigma}^\nu - \sigma^\nu\bar{\sigma}^\mu & 0 \\ 0 & \bar{\sigma}^\mu\sigma^\nu - \bar{\sigma}^\nu\sigma^\mu \end{pmatrix}$$

The Pauli matrices (and the identity matrix) are Hermitian, so:

$$(S^{\mu\nu})^\dagger = -\frac{i}{4} \begin{pmatrix} \bar{\sigma}^\nu\sigma^\mu - \bar{\sigma}^\mu\sigma^\nu & 0 \\ 0 & \sigma^\nu\bar{\sigma}^\mu - \sigma^\mu\bar{\sigma}^\nu \end{pmatrix}$$

Thus:

$$(S^{\mu\nu})^\dagger = \frac{i}{4} \begin{pmatrix} \bar{\sigma}^\mu\sigma^\nu - \bar{\sigma}^\nu\sigma^\mu & 0 \\ 0 & \sigma^\mu\bar{\sigma}^\nu - \sigma^\nu\bar{\sigma}^\mu \end{pmatrix}$$

$$\overline{S^{\mu\nu}} = \frac{i}{4} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} \bar{\sigma}^\mu\sigma^\nu - \bar{\sigma}^\nu\sigma^\mu & 0 \\ 0 & \sigma^\mu\bar{\sigma}^\nu - \sigma^\nu\bar{\sigma}^\mu \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$\overline{S^{\mu\nu}} = \frac{i}{4} \begin{pmatrix} \sigma^\mu \bar{\sigma}^\nu - \sigma^\nu \bar{\sigma}^\mu & 0 \\ 0 & \bar{\sigma}^\mu \sigma^\nu - \bar{\sigma}^\nu \sigma^\mu \end{pmatrix}$$

$$\overline{S^{\mu\nu}} = S^{\mu\nu}$$

□

Next for $i\gamma^5$:

$$\overline{i\gamma^5} = (-i) \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} -I & 0 \\ 0 & I \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

Multiplying the matrices gives:

$$\overline{i\gamma^5} = i\gamma^5$$

as expected. Note that the i contributed a minus sign, so it follows that $\overline{\gamma^5} = -\gamma^5$.

□

Next for $\gamma^\mu \gamma_5$:

$$\overline{\gamma^\mu \gamma_5} = \beta \gamma_5^\dagger \gamma^{\mu\dagger} \beta$$

$\beta^2 = I$, so:

$$\overline{\gamma^\mu \gamma_5} = \beta \gamma_5^\dagger \beta \gamma^{\mu\dagger} \beta$$

We showed above that $\overline{\gamma^\mu} = \gamma^\mu$, and that $\overline{\gamma^5} = -\gamma^5$. Thus:

$$\overline{\gamma^\mu \gamma_5} = -\gamma_5 \gamma^\mu$$

Now, note that $\gamma_5 = i\gamma^0 \gamma^1 \gamma^2 \gamma^3$. The γ^μ matrix will commute with itself but anticommute with the other three. Thus $\{\gamma^\mu, \gamma^5\} = 0$, and:

$$\overline{\gamma^\mu \gamma_5} = \gamma^\mu \gamma_5$$

as expected.

□

Now for $\overline{i\gamma_5 S^{\mu\nu}}$. We have:

$$\overline{i\gamma_5 S^{\mu\nu}} = -i\beta \gamma_5 \beta \beta S^{\mu\nu} \beta$$

where the two β s in the middle are allowed since $\beta^2 = I$. Then,

$$\overline{i\gamma_5 S^{\mu\nu}} = -i\beta \gamma_5 \beta \beta S^{\mu\nu} \beta$$

which is (remember that the gamma matrices including γ_5 are Hermitian, so we can add or remove the dagger at will):

$$\overline{i\gamma_5 S^{\mu\nu}} = -i\overline{\gamma_5} \overline{S^{\mu\nu}}$$

The γ_5 introduces a negative sign as discussed above. We also showed above that $S^{\mu\nu}$ is unchanged under this operation. This gives:

$$\overline{i\gamma_5 S^{\mu\nu}} = i\gamma_5 S^{\mu\nu}$$

□

Srednicki 38.3. Verify equation 38.22.

Subtracting 38.20 from 38.19:

$$\gamma^\mu \not{p} - \not{p}' \gamma^\mu = -p^\mu - 2iS^{\mu\nu} p_\nu + p'^\mu - 2iS^{\mu\nu} p'_\nu$$

We sandwich both sides between $\bar{u}_{s'}(p)$ and $v_s(-p)$ (we will neglect the spinor indices until the end; they are irrelevant to the problem):

$$\bar{u}(p) [\gamma^\mu \not{p} - \not{p}' \gamma^\mu] v(-p) = \bar{u}(p) [-p^\mu - 2iS^{\mu\nu} p_\nu + p'^\mu - 2iS^{\mu\nu} p'_\nu] v(-p)$$

Using equation 38.1:

$$\bar{u}(p) [-m\gamma^\mu - \not{p}' \gamma^\mu] v(-p) = \bar{u}(p) [-p^\mu - 2iS^{\mu\nu} p_\nu + p'^\mu - 2iS^{\mu\nu} p'_\nu] v(-p)$$

Next note that equation 38.16 gives:

$$\begin{aligned} \not{p}v(p) &= mv(p) \\ \implies (-\not{p})v(-p) &= mv(-p) \\ \implies \not{p}v(-p) &= -mv(-p) \end{aligned} \tag{38.3.1}$$

Hence,

$$\bar{u}(p) [-m\gamma^\mu + m\gamma^\mu] v(-p) = \bar{u}(p) [-p^\mu - 2iS^{\mu\nu} p_\nu + p'^\mu - 2iS^{\mu\nu} p'_\nu] v(-p)$$

This gives:

$$\bar{u}(p) [-p^\mu - 2iS^{\mu\nu} p_\nu + p'^\mu - 2iS^{\mu\nu} p'_\nu] v(-p) = 0$$

Next we take $p = p'$:

$$\bar{u}(p) [-4iS^{\mu\nu} p_\nu] v(-p) = 0$$

Dropping the constants:

$$\bar{u}(p) [S^{\mu\nu} p_\nu] v(-p) = 0$$

Using 36.53:

$$\bar{u}(p) \left[\begin{pmatrix} -\sigma^\nu & 0 \\ 0 & \sigma^\nu \end{pmatrix} p_\nu \right] v(-p) = 0$$

We can write this as:

$$\bar{u}(p) \left[\begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^\nu \\ -\sigma^\nu & 0 \end{pmatrix} p_\nu \right] v(-p) = 0$$

These are gamma matrices:

$$\bar{u}(p) [\gamma^0 \gamma^\nu p_\nu] v(-p) = 0$$

Using equation 37.24:

$$\bar{u}(p) \gamma^0 \not{p} v(-p) = 0$$

Using equation (38.3.1):

$$-\bar{u}(p) \gamma^0 mv(-p) = 0$$

Dropping the constants:

$$\bar{u}_{s'}(p)\gamma^0 v_s(-p) = 0$$

□

Taking the Hermitian conjugate:

$$[\bar{u}(p)\gamma^0 v(-p)]^\dagger = 0$$

This gives:

$$v(-p)^\dagger \gamma^{0\dagger} \bar{u}(p)^\dagger = 0$$

γ^0 is Hermitian. On the other terms, we use 38.7 (multiplying on the right by β). The result is:

$$\bar{v}(-p)\beta\gamma^0\beta u(p) = 0$$

Recall that β and γ^0 are the same (we keep them separate only to maintain the integrity of the indices). It is therefore obvious that they commute. Thus

$$\bar{v}(-p)\beta\beta\gamma^0 u(p) = 0$$

$\beta^2 = I$, so, reinserting the indices:

$$\bar{v}_s(-p)\gamma^0 u_{s'}(p) = 0$$

Finally, we let $p \rightarrow -p$ and $s \leftrightarrow s'$:

$$\bar{v}_{s'}(p)\gamma^0 u_s(-p) = 0$$

□

Srednicki 38.4. Derive the Gordon identities:

$$\bar{u}_{s'}(\mathbf{p}') [(\mathbf{p}' + \mathbf{p})^\mu - 2iS^{\mu\nu}(\mathbf{p}' - \mathbf{p})_\nu] \gamma_5 u_s(\mathbf{p}) = 0$$

$$\bar{v}_{s'}(\mathbf{p}') [(\mathbf{p}' + \mathbf{p})^\mu - 2iS^{\mu\nu}(\mathbf{p}' - \mathbf{p})_\nu] \gamma_5 v_s(\mathbf{p}) = 0$$

The technique here is the same as in the previous problem. We start by adding 38.19 and 38.20:

$$\gamma^\mu \not{p} + \not{p}' \gamma^\mu = -p^\mu - p'^\mu - 2iS^{\mu\nu} p_\nu + 2iS^{\mu\nu} p'_\nu$$

We multiply on the right by γ_5 , then sandwich between $\bar{u}(p')$ and $u(p)$. Thus:

$$\bar{u}(p') [\gamma^\mu \not{p} + \not{p}' \gamma^\mu] \gamma_5 u(p) = \bar{u}(p') [-p^\mu - p'^\mu - 2iS^{\mu\nu} p_\nu + 2iS^{\mu\nu} p'_\nu] \gamma_5 u(p)$$

Using equation 38.16:

$$\bar{u}(p') [\gamma^\mu \not{p} - m\gamma^\mu] \gamma_5 u(p) = -\bar{u}(p') [p^\mu + p'^\mu + 2iS^{\mu\nu}(p_\nu - p'_\nu)] \gamma_5 u(p)$$

Recall that γ^5 consists of four gamma matrices: γ^μ will anticommute with three of them and commute with one (itself):

$$\bar{u}(p') [-\gamma^\mu \gamma_5 \not{p} - m\gamma^\mu \gamma_5] \gamma_5 u(p) = -\bar{u}(p') [p^\mu + p'^\mu + 2iS^{\mu\nu}(p_\nu - p'_\nu)] \gamma_5 u(p)$$

Using 38.1:

$$\bar{u}(p') [m\gamma^\mu\gamma_5 - m\gamma^\mu\gamma_5] \gamma_5 u(p) = -\bar{u}(p') [p^\mu + p'^\mu + 2iS^{\mu\nu}(p_\nu - p'_\nu)] \gamma_5 u(p)$$

This left side vanishes:

$$-\bar{u}(p') [p^\mu + p'^\mu + 2iS^{\mu\nu}(p_\nu - p'_\nu)] \gamma_5 u(p) = 0$$

Dropping the constant and reordering:

$$\bar{u}(p') [p^\mu + p'^\mu - 2iS^{\mu\nu}(p' - p)_\nu] \gamma_5 u(p) = 0$$

□

For the remaining identity, we repeat this with $u \leftrightarrow v$. We see from 38.1 and 38.16 that the result will be the same up to a minus sign; since there are two minus signs, there is no net result. Thus:

$$\bar{v}(p') [p^\mu + p'^\mu - 2iS^{\mu\nu}(p' - p)_\nu] \gamma_5 v(p) = 0$$

□