# QFT

#### **Chapter 53: Functional Determinants**

#### **Overview**

• Recall in chapter 44, we introduced Grassman numbers to derive equation 44.40, a general mathematical formula:

$$\int d^n \chi d^n \overline{\chi} \exp(\chi^{\dagger} M \chi + \eta^{\dagger} \chi + \chi^{\dagger} \eta) = (\det M) \exp(-\eta^{\dagger} M^{-1} \eta)$$

- The fermionic path integral was just a special case of this.
- In this special case, the determinant was independent of the fields and sources, and could just be absorbed into the overall normalization.
- Here, we will ponder the meaning of this will be needed later on when we get to non-Abelian Gauge theories.

### Analysis #1

We have from chapter 44 that

$$\int d^n z d^n \overline{z} e^{-i\overline{z}_i M_{ij} z_j} \propto (\det M)^{-1}$$

• In the case at hand, we have:

$$M = \left[-\partial_x^2 + m^2 - g\phi\right]\delta^4(x - y)$$

• But how do we take the determinant of this? First, we define:

$$M = \int d^4 y M_0(x, y) \widetilde{M}(y, z)$$

- This factorization is non-intuitive, but Srednicki "guesses" the proper factorization and it is easy to verify that he is right
- The advantage is that  $M_0$  is defined to be independent of the background field  $\phi,$  and can therefore be absorbed into the normalization

#### Analysis #1, continued

• This gives us:

$$Z(\phi) = (\det \widetilde{M})^{-1}$$

• As it turns out, the form needed to factor was:

$$\widetilde{M}(y,z) = \delta^4(y-z) - g\Delta(y-z)\phi(z)$$

• So now we need to take the determinant of this. We can use the general matrix relation

$$\det A = \exp \operatorname{Tr} \log A$$

• and it follows that we get:

$$Z(\phi) = \exp\sum_{n=1}^{\infty} \frac{1}{n} \left[ g^n \int d^4x_1 \dots d^4x_n \Delta(x_1 - x_2)\phi(x_2) \dots \Delta(x_n - x_1)\phi(x_1) \right]^n$$

## Analysis #2

- Now we can redo this, treating the interaction term in the Lagrangian as an interaction (and phi as part of the vertex factor, rather than a separate field).
  - The only diagrams we can draw are circles of vertices.
  - Summing the diagrams, we rederive our result from analysis #1

# Again

- Now we repeat both these analyses for fermions.
  - The difference is the negative sign we found in problem 51.1

 It's important to note that not only did we derive the negative sign, but we also showed that we can deal directly with these functional determinants.