

QFT

Chapter 30: Spontaneous Symmetry Breaking

ϕ^4 Theory, again

- Consider ϕ^4 theory:

$$\mathcal{L} = -\frac{1}{2}\partial^\mu\phi\partial_\mu\phi - \frac{1}{2}m^2\phi^2 - \frac{1}{24}\lambda\phi^4$$

- Some obvious observations:

- We can swap $\phi \rightarrow -\phi$
- In operator form, we can define $Z^{-1}\phi Z = -\phi$
- Then $Z^2 = 1$, which implies $Z = Z^{-1}$
- By definition, Z is unitary, so $Z^{-1} = Z^\dagger$
- This implies $Z = Z^\dagger$, ie Z is Hermitian

- We can write this as:

$$\mathcal{L} = -\frac{1}{2}\partial^\mu\phi\partial_\mu\phi - \frac{1}{24}\lambda(\phi^2 - v^2)^2 - \frac{1}{24}\lambda v^4$$

where $v = \sqrt{6|m^2|/\lambda}$. The last term is constant (at least, it doesn't have any ϕ s), so we drop it.

Ground State

$$\mathcal{L} = -\frac{1}{2}\partial^\mu\phi\partial_\mu\phi - \frac{1}{2}m^2\phi^2 - \frac{1}{24}\lambda\phi^4$$

$$\mathcal{L} = -\frac{1}{2}\partial^\mu\phi\partial_\mu\phi - \frac{1}{24}\lambda(\phi^2 - v^2)^2$$

- What is the ground state of this Lagrangian?
 - If $m^2 > 0$, then we can easily see from the original Lagrangian that there is only one ground state at $\phi = 0$. Then $\langle 0|\phi|0\rangle = 0$, just as required by the renormalization scheme.
 - If $m^2 < 0$, then we see from the new Lagrangian that the ground is clear: $\phi = \pm v$.
- In QFT also (with $m^2 < 0$), we expect two ground states, 0_+ and 0_- . These are exchanged by the operator ϕ , and are orthogonal.
 - Why orthogonal? Consider space as a collection of harmonic oscillators. You can tunnel from 0_+ to 0_- , but all infinitely many would have to do so, which lowers the amplitude to 0.

Broken Symmetry

- Let's imagine that 0^+ is the ground state (when $m < 0$). Let's also define $\rho = \phi - v$. This is required to fix our renormalization condition $\langle 0|\rho|0\rangle$.

$$\mathcal{L} = -\frac{1}{2}\partial^\mu\phi\partial_\mu\phi - \frac{1}{6}\lambda v^2\rho^2 - \frac{1}{6}\lambda v\rho^3 - \frac{1}{24}\lambda\rho^4$$

- The second term coefficient could be set to $\frac{1}{2}m_\rho^2$, the mass of the ρ particle.
- But there is a problem: where the hell is our Z_2 symmetry?
 - Spontaneously broken!

Symmetry Breaking in the Quantum Action

- Recall in problem 21.2, we showed that the quantum action inherits any linear symmetry of the classical action, provided that it is also a symmetry of the integration measure (which it usually is).
- The quantum action “therefore” shows the spontaneous symmetry breaking as well.
- All the ingredients of $\Gamma(\varphi)$ are made finite with the three existing Z s, so no additional Z s are needed.
- We’ll see how this works in the next section.