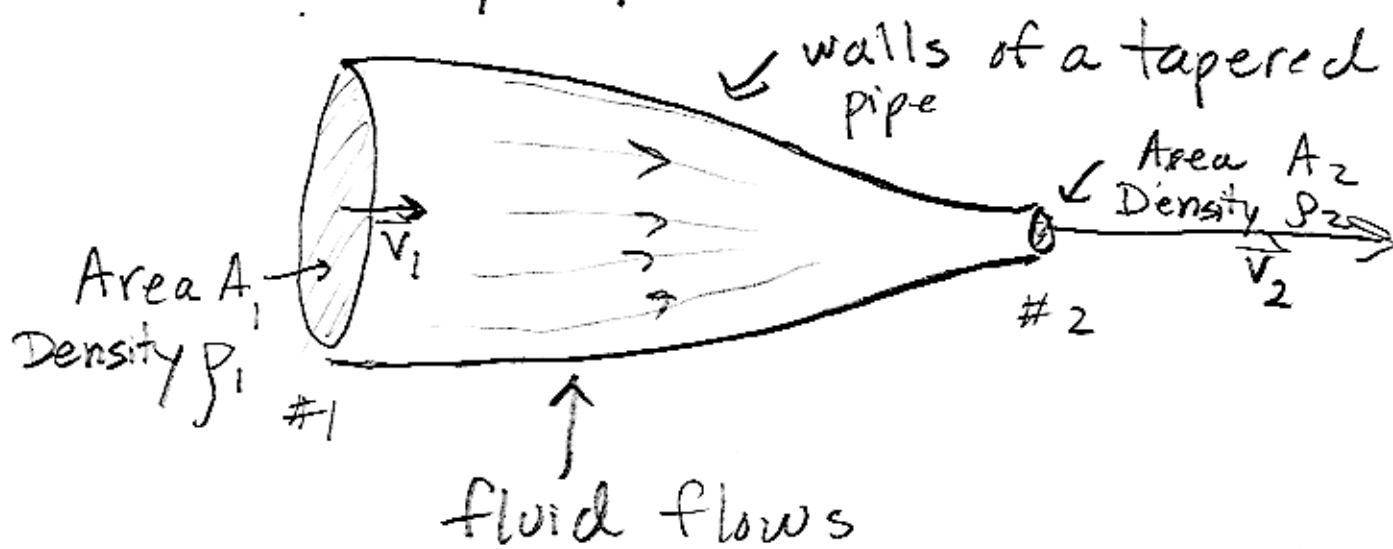


Chapter 23: (Gauss's Law)

1. The Main Idea
2. Electric Flux
3. Gauss's Law
4. Applications

The Main Idea - develop an analogy to the continuity equation (which is for fluids)... analogy is for electric field.

Continuity Equation was:



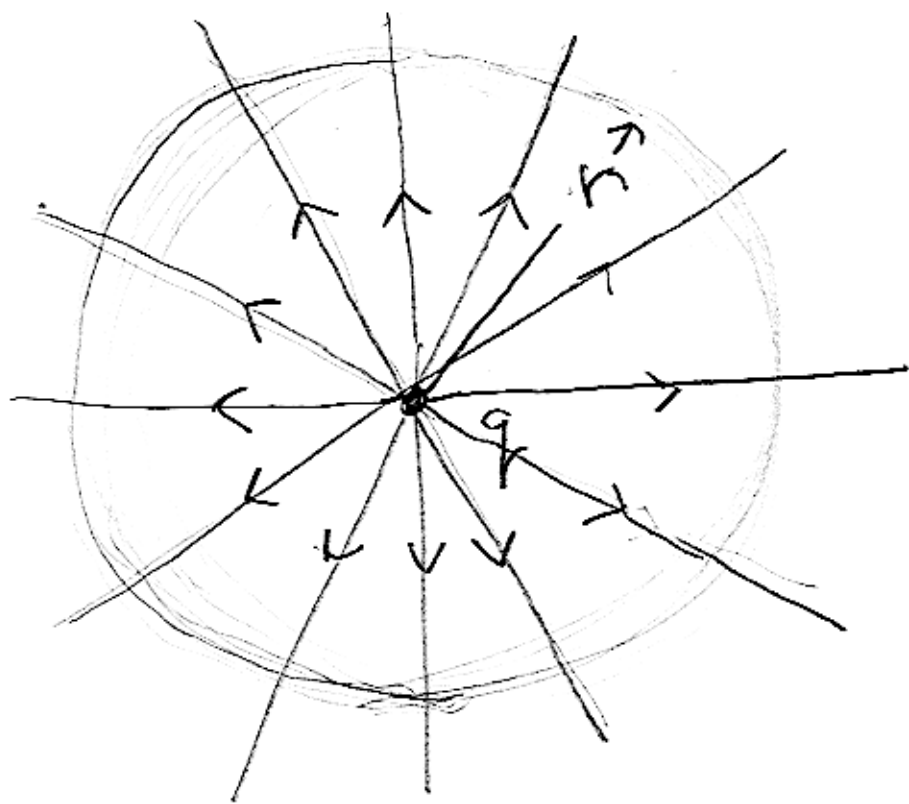
in time dt , same mass flowing in through area A_1 , must flow out through surface A_2 :

$$(\rho_1)(A_1 v_1 dt) = (\rho_2)(A_2 v_2 dt)$$

$$\rho_1 A_1 v_1 = \rho_2 A_2 v_2$$

A similar thing happens for the correctly used quantity involving the electric field.

To develop that quantity, imagine 59
the electric field surrounding a + point
charge: Imagine a spherical surface
centered on the point charge too:



The radius of the surface is r

On the surface of the sphere, the
magnitude of electric field is:

$$E_{\text{surface}} = k \frac{q}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}$$

Now note that the surface area of
the sphere is $4\pi r^2$, so,

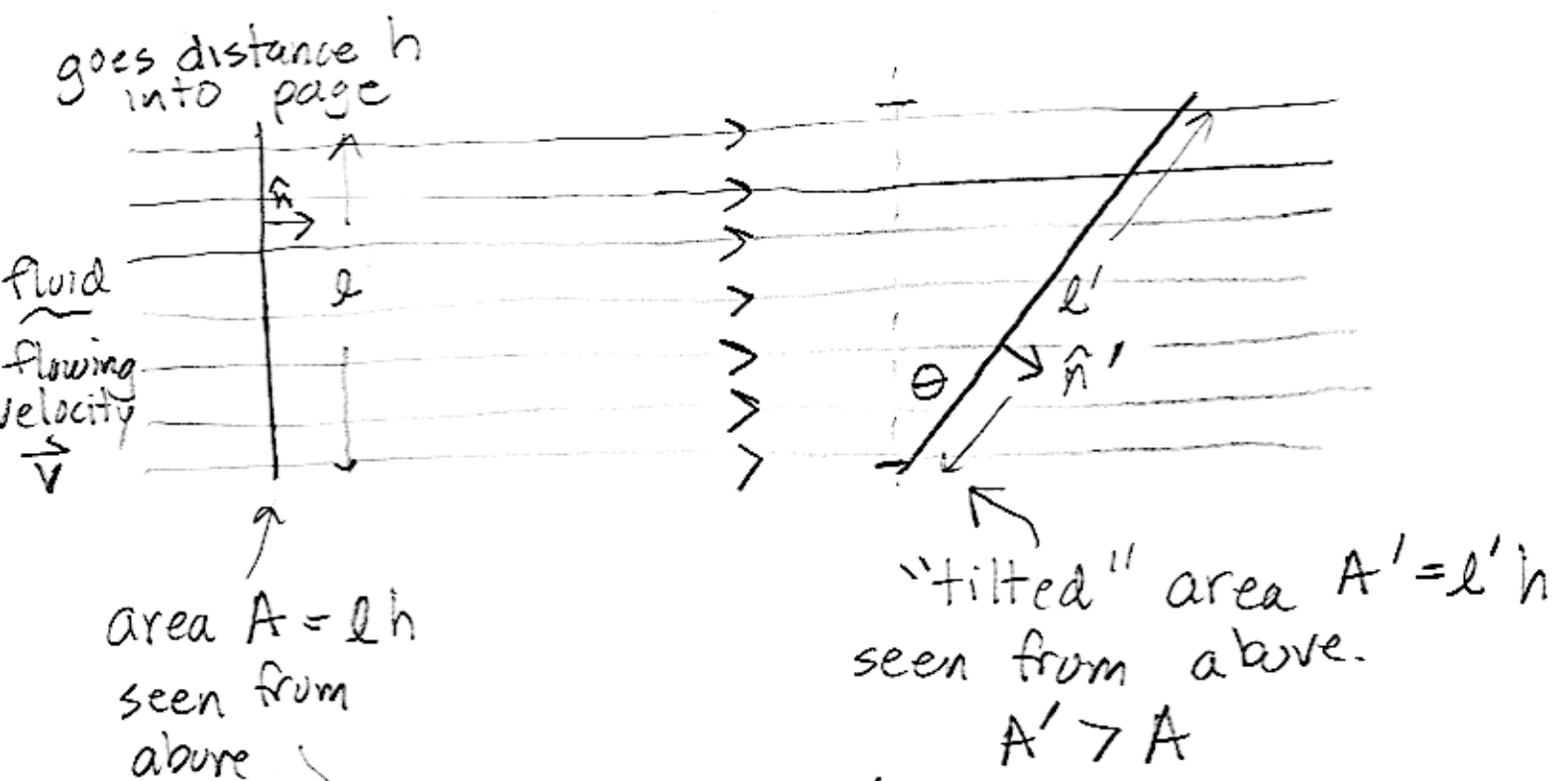
$$\underbrace{(\text{Surface Area}) \times E_{\text{surface}}}_{\text{"electric flux"}} = 4\pi r^2 \times \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r^2} = \frac{q}{\epsilon_0} \quad (\text{Independent of Radius of surface!})$$

To further the analogy, identity: 60

$$\vec{E} \iff \rho \vec{v}$$

and note that mass flow (described by $\rho \vec{v}$) is the same even when the surface it flows through is tilted w/r to \vec{v} :

Look down from above. ρ, \vec{v} constant:



$$\rho v A \neq \rho v A'$$

but: $\rho v A = \rho v A \cos \theta$ since $l = l' \cos \theta$

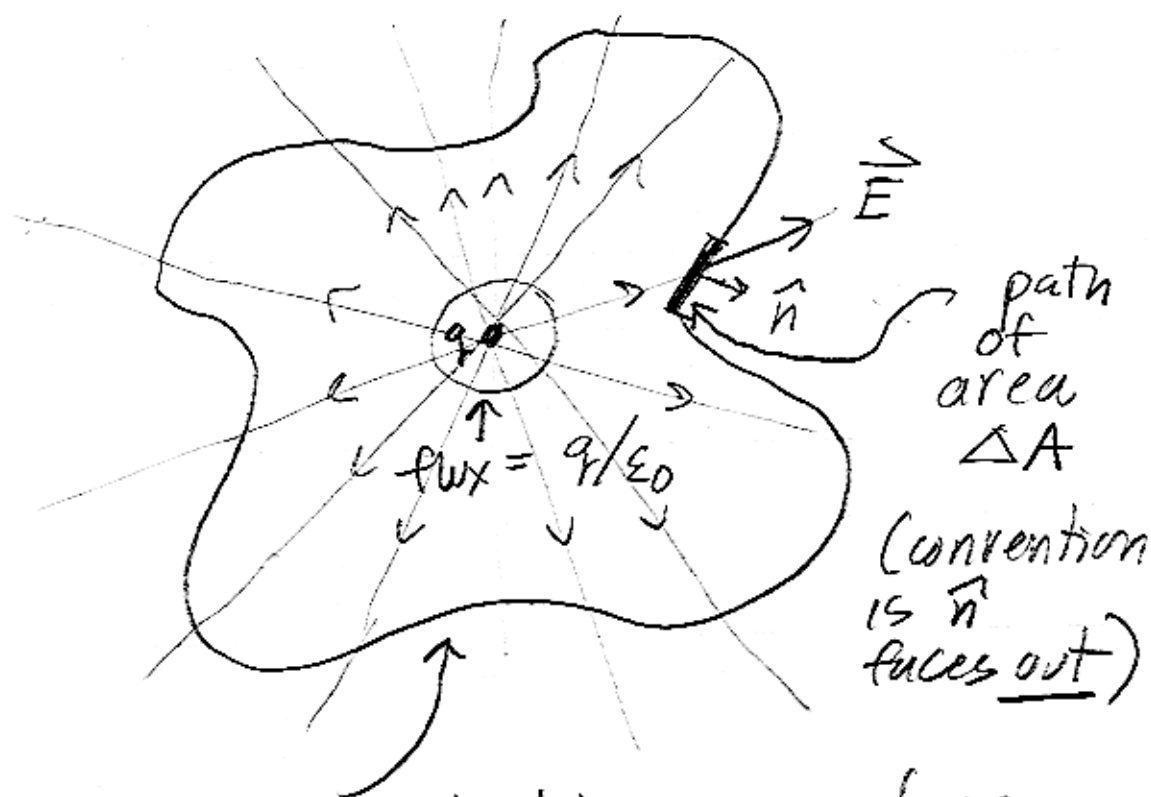
$$= \rho (v \cos \theta) A$$

\vec{v} 's component \perp to surface.

$\hat{n} \equiv$ unit vector normal to surface

then $\rho A \vec{v} \cdot \hat{n} = \rho A' \vec{v} \cdot \hat{n}'$ "mass flux"

Now (try) to imagine situation for electric field... idea is that the electric flux through a small sphere that surrounds a point charge will also "flow" through an irregular surface that also fully surrounds the charge:



flux through bigger surface (closed, enclosing charge) ALSO $\frac{q}{\epsilon_0}$

How to define the flux for the larger, irregular surface?

$$\begin{aligned}\Phi_E &\equiv \sum \vec{E} \cdot \hat{n} \Delta A \Rightarrow \oint \vec{E} \cdot \hat{n} dA \\ &= \oint \vec{E} \cdot d\vec{A} \quad d\vec{A} \equiv \hat{n} \\ &= 4\pi r^2 \cdot \frac{q}{4\pi \epsilon_0 r^2} = \frac{q}{\epsilon_0} \leftarrow\end{aligned}$$

Back to nice spherical surface

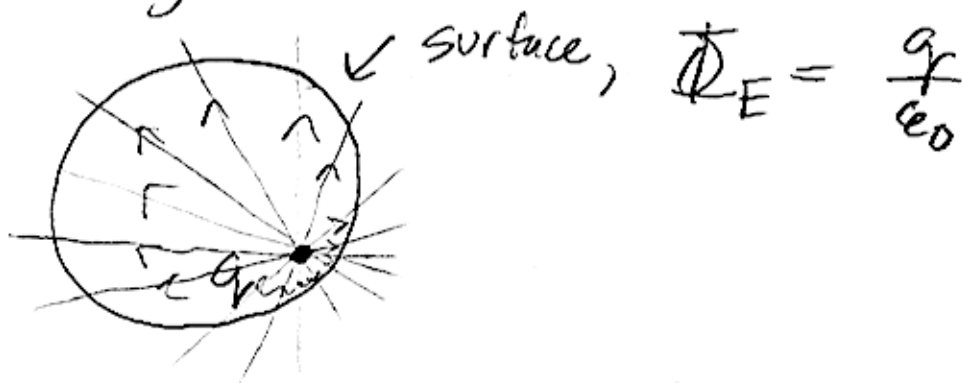
note that if $E(r) \neq \frac{1}{r^2}$ for point charge, flux would vary as a function of the size of enclosing surface.
 Since $E(r)$ is $\propto 1/r^2$,

$$\Phi_E = \oint \vec{E} \cdot d\vec{A} = \frac{Q_{\text{enclosed}}}{\epsilon_0}$$

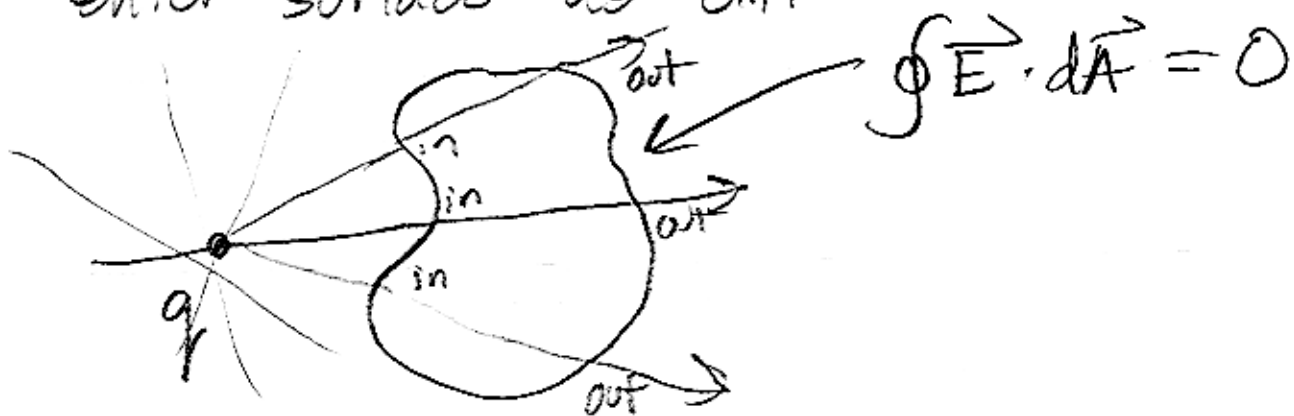
↑
means over
a closed surface

Gauss's Law

1) Charge need not be in center!

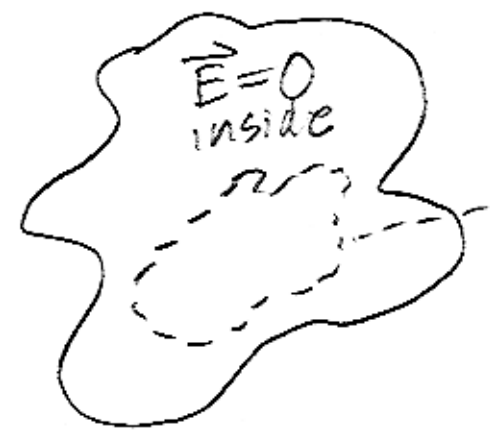


2) When all charge is outside, the net flux is 0... equal amounts enter surface as exit.



3) Static charge on an ideal piece of a conductor resides on surface.

⇒ key insight: no electric field inside an ideal conductor (or else current would flow, pushed by field) at static equilibrium.



↑
boundary of
block of ideal
conductor

→ arbitrary surface fully inside conductor

→ $\Phi_E = 0$ because

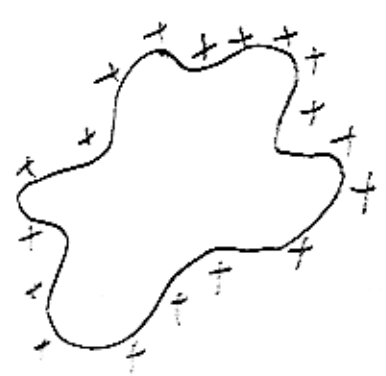
$\vec{E} = 0$ everywhere inside.

→ ∴ $Q_{\text{enclosed}} = 0$ for arbitrary surface.

→ ∴ no charge ever inside conductor.

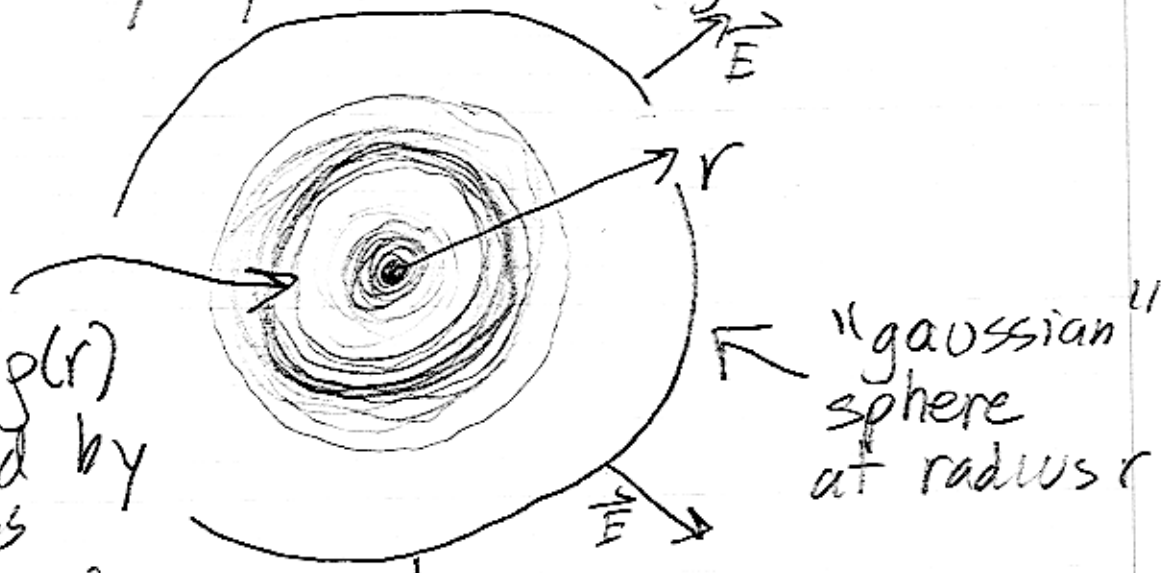
→ but you can charge up a conductor!

so all charge lives on the surface!



4) Spherically symmetric charge distributions:

3-d charge density $\rho(r)$ indicated by darkness (function of r only)



\vec{E} always radially outward (or inward)

Flux
 $4\pi r^2 E_r$

$$= \frac{\text{(Charge Enclosed)}}{\epsilon_0}$$

radial component of r

$$E_r = \frac{\text{(Charge Enclosed)}}{4\pi\epsilon_0 r^2}$$

This means that for small r , the charge OUTSIDE r has no effect on E_r !



← this charge doesn't influence E_r

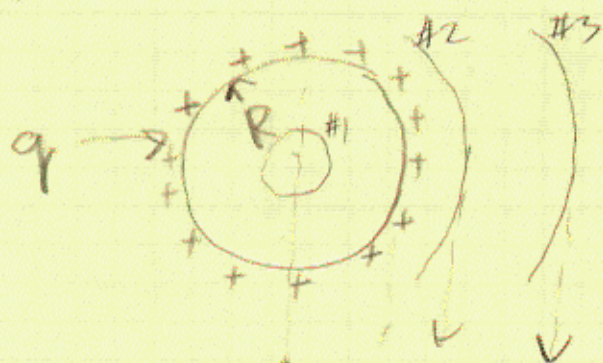
5) Hole inside a spatial charge distribution

~~WRONGS!~~

darkness indicates $\rho(x)$

Hole: $\vec{E} = 0$ inside, because every gaussian surface inside encloses 0 charge, so $\Phi_E = 0$ for arbitrary surface, $\vec{E} = 0$

6) Field of a charged conducting sphere:

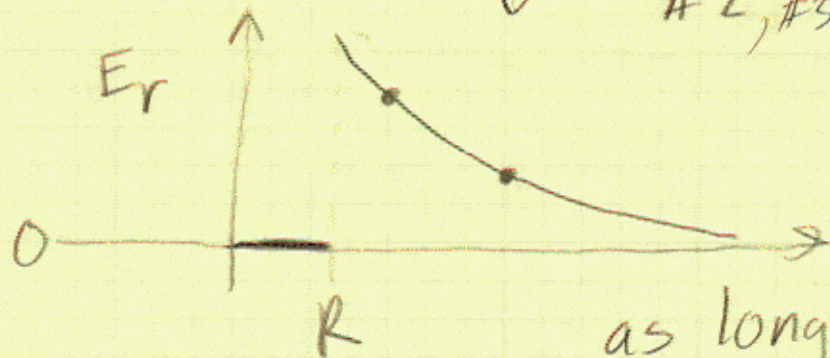


Gaussian Surfaces (portions shown)

#1: no enclosed charge
 $E_r = 0 \dots$ up to $r = R$

#2, #3: $4\pi r^2 E_r = \frac{q}{\epsilon_0}$

$$E_r = \frac{q}{4\pi\epsilon_0 r^2}$$



as long as $r > R$,
field just like a
point charge!

(Newton knew this too, but he didn't know Gauss's law.)