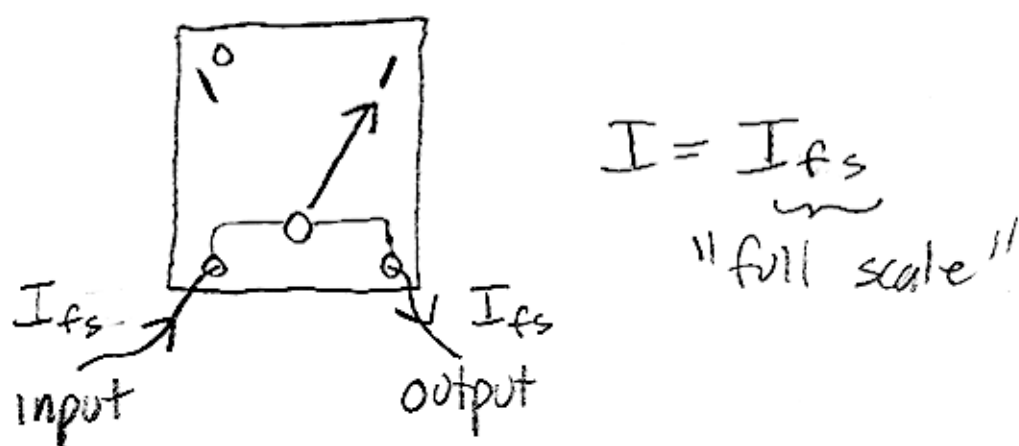
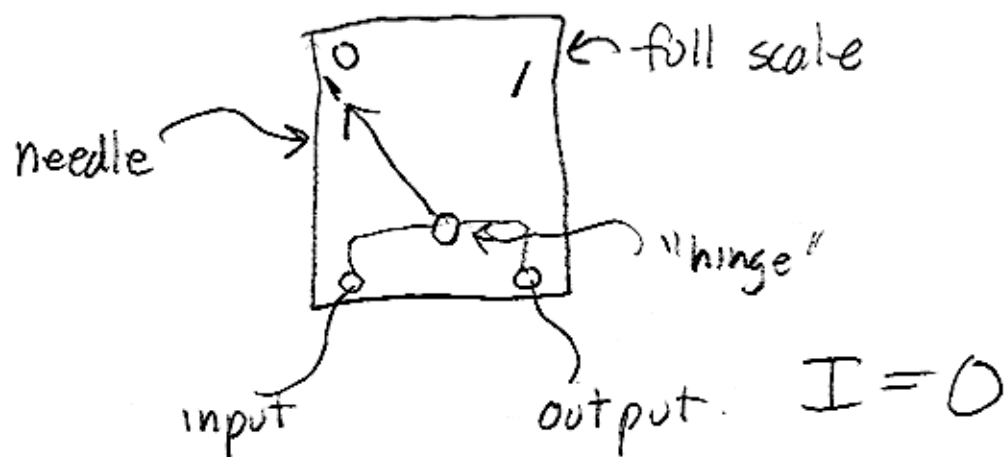


Old-Style Meters

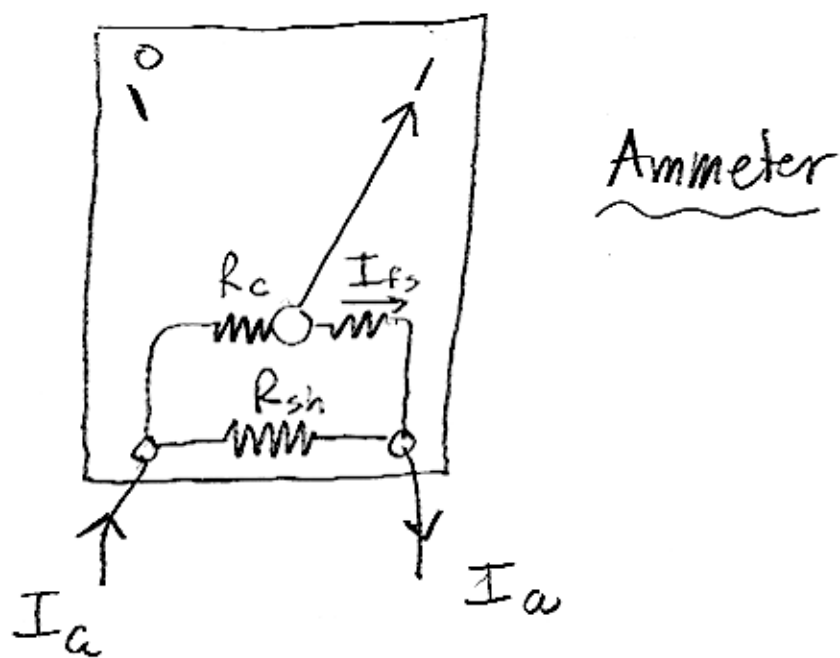
Are usually "galvanometers" which convert electrical current into angular deflection of a needle.



- 1) Physical mechanism is: current causes a magnetic field, magnetic field causes a torque.
- 2) Field caused by a "coil", which actually has a non-zero resistance R_c ... gives non-ideal behavior; $R_c \approx 0 \Omega$
- 3) Art is to use parallel and series resistances to make full scale deflection correspond to desired behavior.

Example... make full scale correspond to less than I_{fs} by "bleeding" current off through a "shunt" resistor, R_{sh}

goal: make $I_a > I_{fs}$ give the full scale reading



$$R_c I_{fs} = R_{sh} (I_a - I_{fs})$$

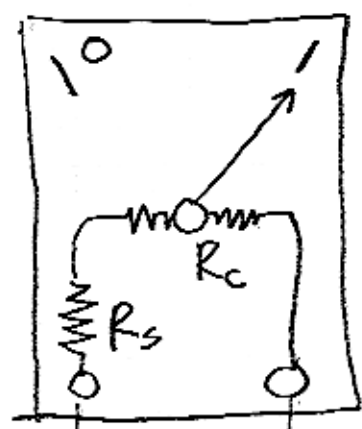
$$R_{sh} = \frac{I_{fs}}{I_a - I_{fs}} \cdot R_c, \quad > R_c \quad \text{when } I_a > I_{fs}$$

$$\frac{1}{R_{eq}} = \frac{1}{R_c} + \frac{1}{R_{sh}}$$

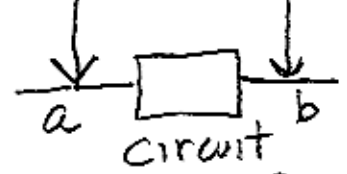
gives equivalent resistance, $< R_c$ but not 0.

Voltmeter

put resistor R_s
in series



$$V_{abmax} \equiv V_V$$



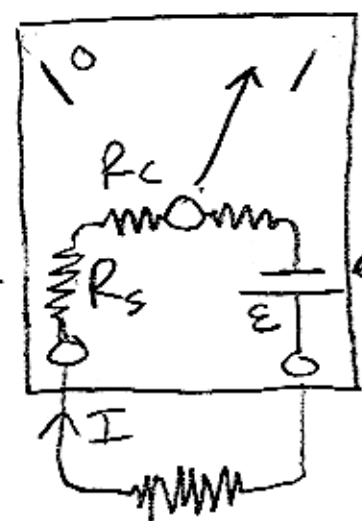
need R_{eq} for circuit
 $\ll (R_s + R_c)$

$$\frac{V_V}{R_s + R_c} = I_{fs}$$

$$R_s = \frac{V_V}{I_{fs}} - R_c, \text{ usually } \gg R_c$$

Ohmmeter

pick this \rightarrow



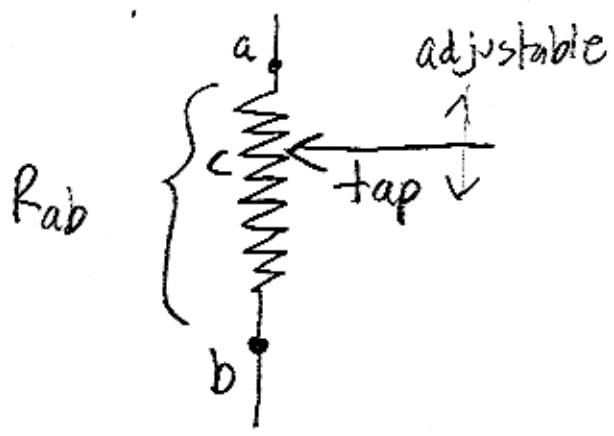
new battery introduced

$R \leftarrow$ measure this
(scale looks funny)

$$I = \frac{E}{R + R_s + R_c}$$

Potentiometer : Measure the voltage of an unknown battery

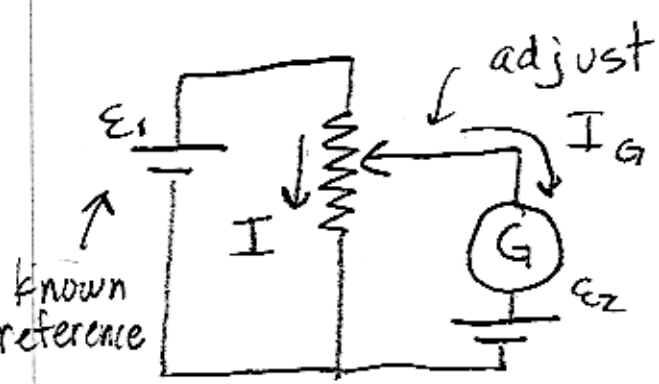
↓ resistor with a "tap"



$$R_{ac} < R_{ab}$$

$$R_{cb} = R_{ab} - R_{ac}$$

Use this, with a galvanometer, to make more accurate measurement of the voltage across a battery.



$$I R_{cb} = \epsilon_2$$

$$I \approx \frac{\epsilon_1}{R_{ab}} \text{ (or } \frac{\epsilon_1}{R_{ab} + r} \text{)}$$

$$\text{so } \epsilon_2 = \frac{R_{cb}}{R_{ab}} \cdot \epsilon_1$$

only works for $\epsilon_2 < \epsilon_1$

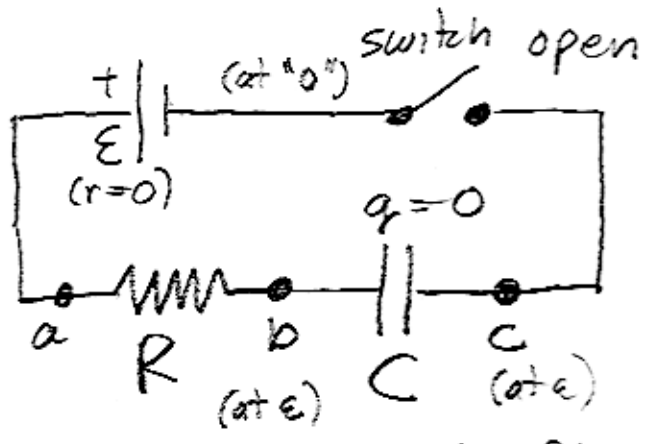
o No current need flow in/cot of ϵ_2 battery ...

Resistance-Capacitance Circuits

V, I, Q : capital letters mean constant in time.

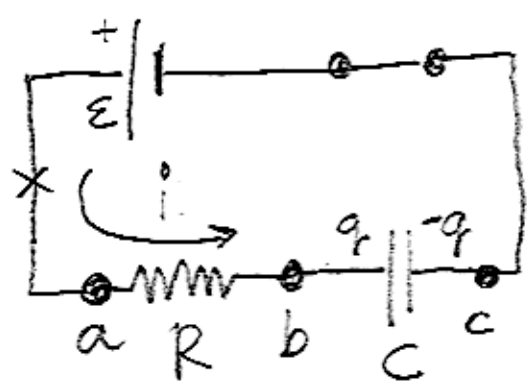
v, i, q : small letters mean varying in time (possibly)

Initial :
from $t = -\infty$
up to $t = 0$



switch open: no current flows.
capacitor uncharged: $q = 0$
 $V_{ab} = V_{bc} = 0$ so all at "ε"

Now close switch at $t = 0$ (this instant in time is singular); analyze voltages, currents, charges as a function of subsequent time t



charge conservation:
 $i = \frac{dq}{dt}$

loop :

$$-iR - \frac{q}{C} + \epsilon = 0$$

$$-R \frac{dq}{dt} - \frac{q}{C} + \epsilon = 0$$

Qualitative Reasoning

① $t=0$, $q(0)=0$ so $-R \left. \frac{dq}{dt} \right|_{t=0} - \frac{0}{C} + \mathcal{E} = 0$

$q(0) = 0$	$i(0) = \left. \frac{dq}{dt} \right _{t=0} = \frac{\mathcal{E}}{R}$
------------	---

② $t \Rightarrow \infty$, all voltage drop is across the capacitor now: charge has built up to the point where a static situation has been achieved:

$$i(t=\infty) = 0 = \left. \frac{dq}{dt} \right|_{t=\infty}$$

$$-R \cdot 0 - \frac{q(\infty)}{C} + \mathcal{E} = 0$$

$i(\infty) = 0$	$q(\infty) = C\mathcal{E}$
-----------------	----------------------------

what is "transition time"

Quantitative Solution

(change $q \rightarrow q'$)
just notation

$$-R \frac{dq'}{dt'} - \frac{q'}{C} + \mathcal{E} = 0$$

$$-R \frac{dq'}{dt'} = -\mathcal{E} + \frac{q'}{C}$$

$$\frac{dq'}{dt'} = \frac{\mathcal{E}}{R} - \frac{q'}{RC} = \frac{C\mathcal{E} - q'}{RC}$$

dimensionless $\frac{dq'}{q' - C\mathcal{E}} = - \frac{dt'}{RC}$ (RC has dimensions of time)

$$\int_0^q \frac{dq'}{q' - C\varepsilon} = -\frac{1}{RC} \int_0^t dt'$$

$$\ln(q' - C\varepsilon) \Big|_0^q = -\frac{t}{RC}$$

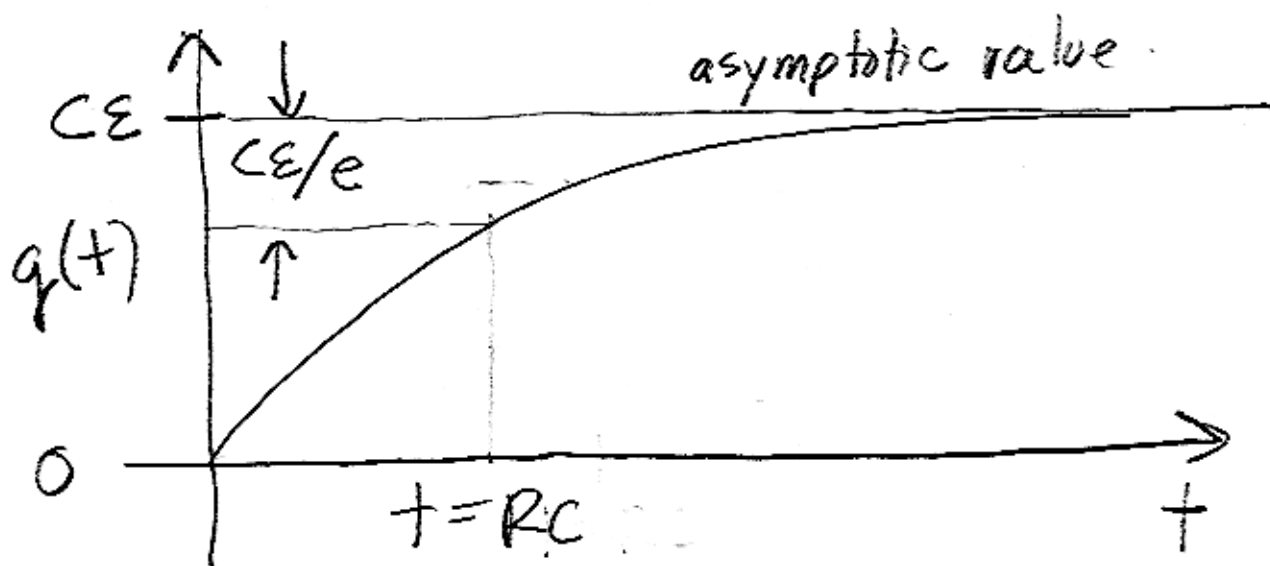
$$\ln\left(\frac{q - C\varepsilon}{-C\varepsilon}\right) = -\frac{t}{RC}$$

$$\frac{q - C\varepsilon}{-C\varepsilon} = e^{-(t/RC)}$$

$$q(t) = C\varepsilon(1 - e^{-(t/RC)})$$

$$q(0) = C\varepsilon(1 - e^{-0}) = C\varepsilon(1 - 1) = 0$$

$$q(\infty) = C\varepsilon(1 - e^{-\infty}) = C\varepsilon$$



$$q(t=RC) = C\varepsilon(1 - e^{-RC/RC}) = C\varepsilon(1 - e^{-1})$$

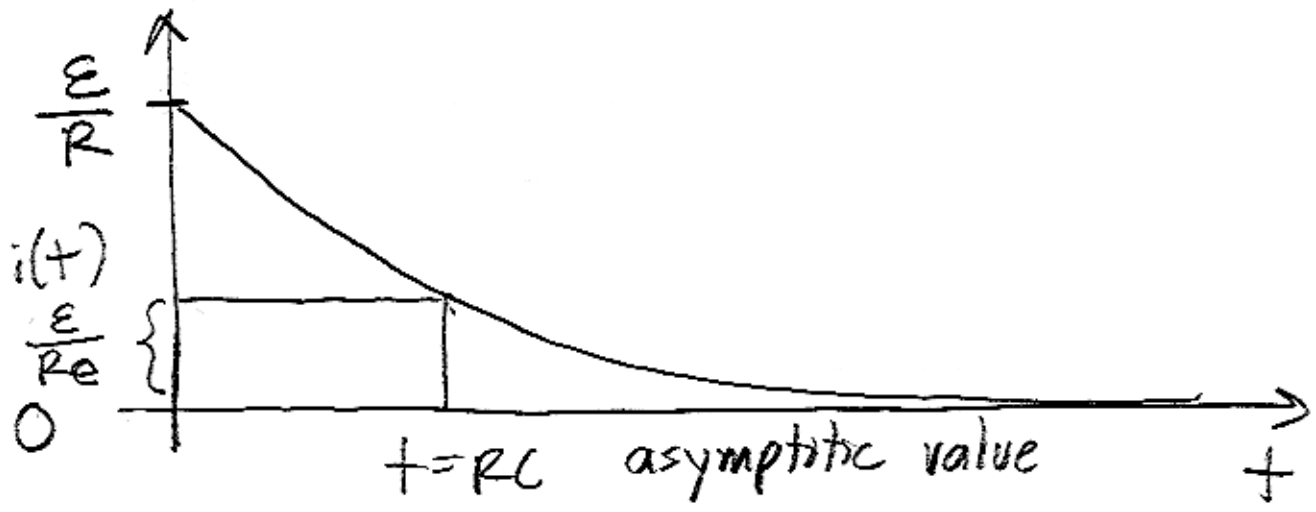
$$= C\varepsilon\left(1 - \frac{1}{e}\right)$$

$$i(t) = \frac{dq}{dt} = \frac{d}{dt} (CE(1 - e^{-t/RC}))$$

$$= -CEe^{-t/RC} \times \left(-\frac{1}{RC}\right) = \frac{\epsilon}{R} e^{-t/RC}$$

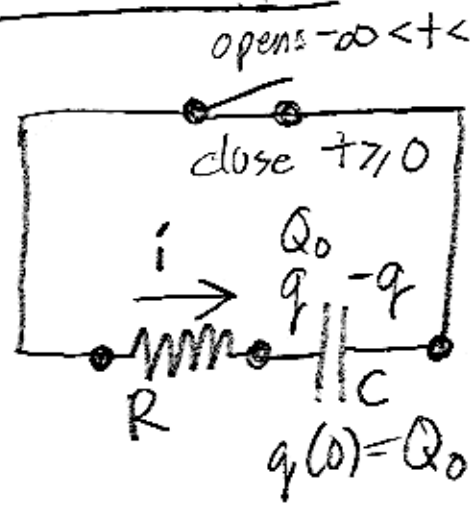
$$i(t) = \frac{\epsilon}{R} e^{-t/RC}$$

$$i(0) = \frac{\epsilon}{R} \quad i(\infty) = 0 \quad i(RC) = \frac{\epsilon}{Re}$$



$\tau = RC$ is the "time constant" of the circuit.

Discharge



$$-iR - \frac{q}{C} = 0$$

$$\frac{dq'}{dt'} = -\frac{q'}{RC}$$

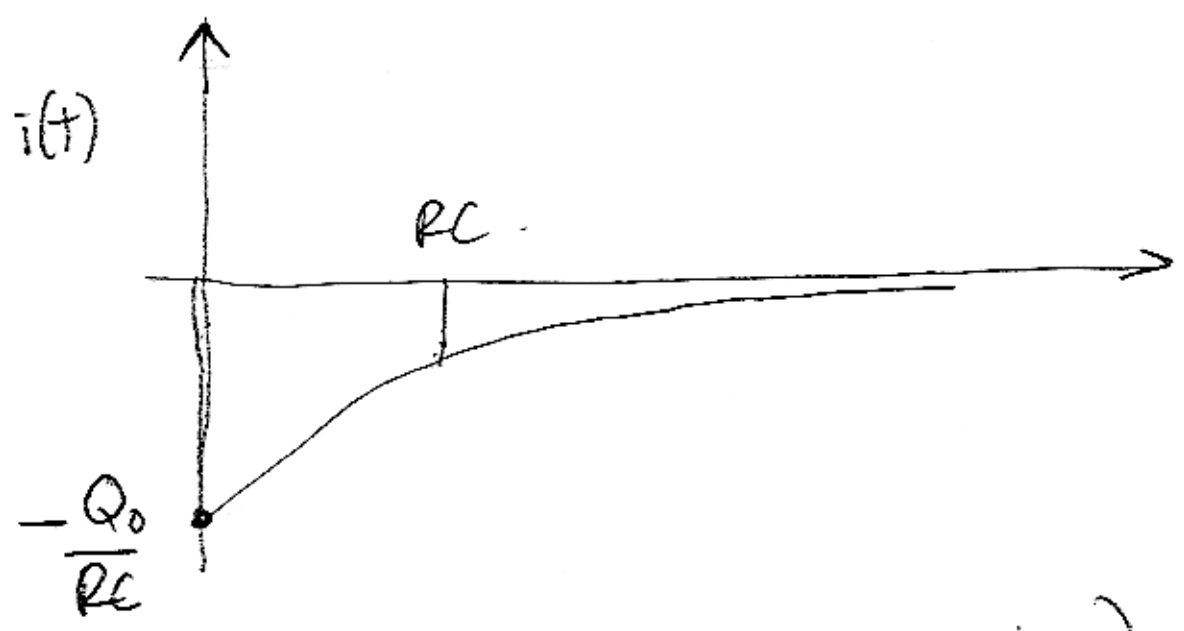
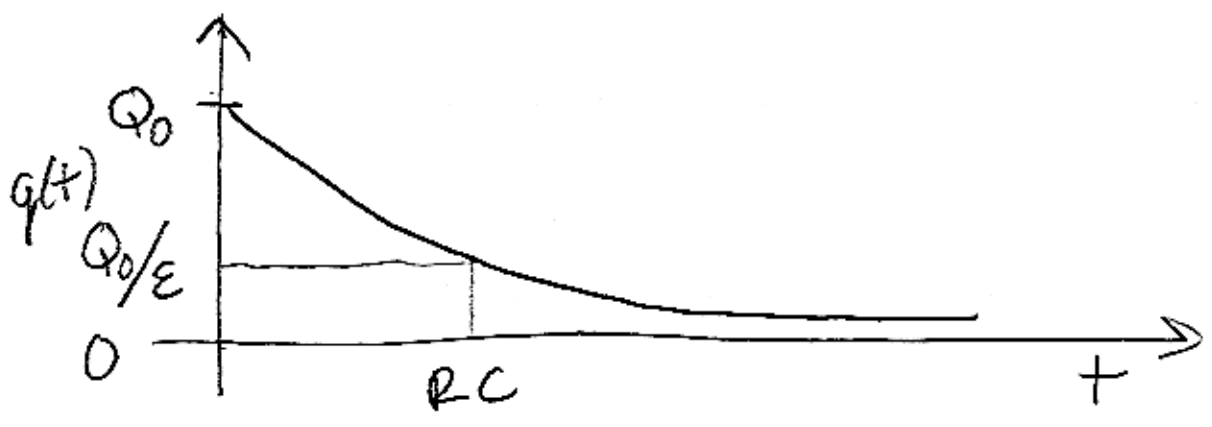
$$\frac{dq'}{q'} = -\frac{dt'}{RC}$$

$$\int_{Q_0}^q \frac{dq'}{q'} = -\frac{1}{RC} \int_0^t dt'$$

$$\ln \frac{q}{Q_0} = -\frac{t}{RC}$$

$$q(t) = Q_0 e^{-t/RC}$$

$$i(t) = \frac{dq}{dt} = -\frac{Q_0}{RC} e^{-t/RC}$$



Energy (Charging up capacitor)

Power Dissipated on Resistor: (charging up)

$$= i^2(t) \cdot R = \frac{\epsilon^2}{R^2} \cdot R \cdot e^{-2t/RC}$$

Energy Lost:

$$\int_0^{\infty} \left(\frac{\epsilon^2}{R}\right) e^{-\frac{2t}{RC}} dt = \frac{\epsilon^2}{R} \left(-\frac{RC}{2}\right) \left(e^{-\frac{2t}{RC}} \Big|_0^{\infty}\right)$$

$$= \frac{1}{2} C \epsilon^2 = \frac{1}{2} \underbrace{(C \epsilon)}_{Q \text{ on capacitor}} \cdot \epsilon$$

Energy Stored on Capacitor

$$= \frac{1}{2} QV = \frac{1}{2} (C \epsilon) \epsilon$$

Energy Provided by Battery

$$= \underbrace{\frac{1}{2} (C \epsilon) \epsilon}_{\text{heats battery}} + \underbrace{\frac{1}{2} (C \epsilon) \epsilon}_{\text{stored on capacitor}}$$

$$= \underline{\underline{C \epsilon^2}} \quad \text{total}$$