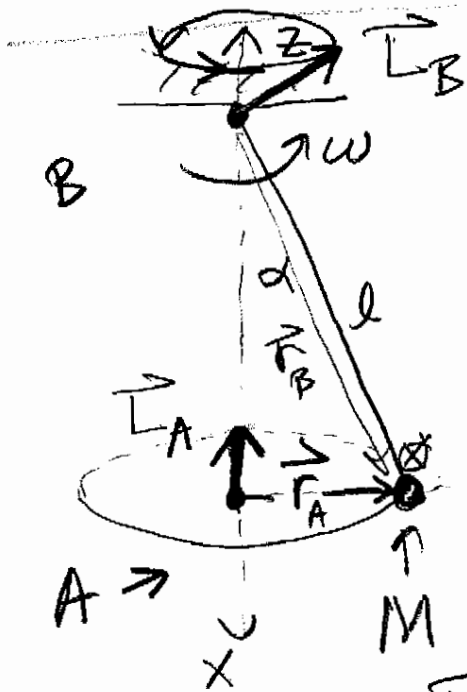


# Conical Pendulum



$$r = l \sin \alpha$$

$$\vec{v} = -r\omega \hat{y}$$

$$\vec{p} = -Mr\omega \hat{y}$$

is  $\perp$  to  $\vec{r}_A$

$$\vec{L}_A = \vec{r}_A \times \vec{p} = Mr^2\omega \hat{k}$$

Constant in magnitude + direction.

note this "moment of inertia"

$$\vec{L}_B = \vec{r}_B \times \vec{p} \quad \vec{r}_B + \vec{p} \text{ also } \perp!$$

$$|\vec{L}_B| = |\vec{r}_\perp| |\vec{p}| = lMr\omega = Mr l \omega$$

magnitude never changes

direction

always in the same plane as mass + suspension point,  $\perp$  to rope.

Note: z-component of  $\vec{L}_B$

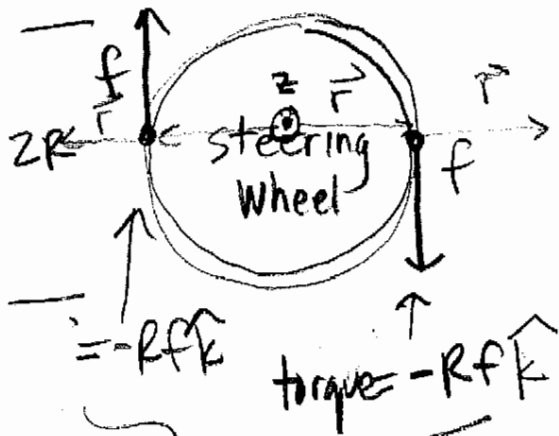
$$= Mr l \omega \sin \alpha = Mr^2 \omega \quad \text{same as } \vec{L}_A$$

Torque

$\vec{\tau} = \vec{r} \times \vec{F}$  ... can decompose  
 just like  $\vec{r} \times \vec{p}$

$|\vec{\tau}| \propto |\vec{r}|$  ... the bigger  $|\vec{r}|$ , the more torque.  
 $\propto \sin(\text{angle between})$   
 push  $\perp$ !

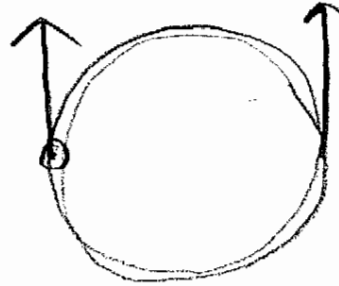
Not the same as force.



add up ...

Net Torque ...  $-2fR\hat{k}$

Net Force ...  $0$



$= 0$

$= 2f\hat{j}$



$= +fR\hat{k}$

$= f\hat{j}$

Torque  $\Leftrightarrow$  Angular Momentum

$\frac{d}{dt}(\vec{L}) = \frac{d}{dt}(\vec{r} \times \vec{p}) \Leftarrow$  chain rule works!

$$\frac{d\vec{L}}{dt} = \frac{d\vec{r}}{dt} \times \vec{p} + \vec{r} \times \frac{d\vec{p}}{dt}$$

$$= \underbrace{\vec{v} \times m\vec{v}}_0 + \vec{r} \times \vec{F}_{\text{Net}} \quad \left. \begin{matrix} \vec{F}_{\text{Net}} = \frac{d\vec{p}}{dt} \end{matrix} \right\}$$

$$\frac{d\vec{L}}{dt} = \vec{r} \times \vec{F}_{\text{Net}}$$

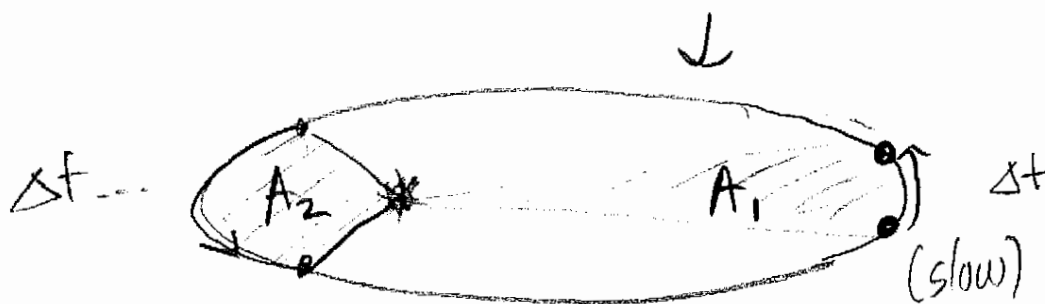
When  $\vec{F}_{\text{Net}} \propto \vec{r}$ ,  $\frac{d\vec{L}}{dt} = 0!$

$\vec{L} = \text{constant}$  (ponder  
critical pendulum)

Most Famous Case... planets about  
sun...

Leads to Kepler's 2<sup>ND</sup> Law...

In equal times, planet's orbits  
sweep out equal area.



In same  $\Delta t$   $A_1 = A_2$