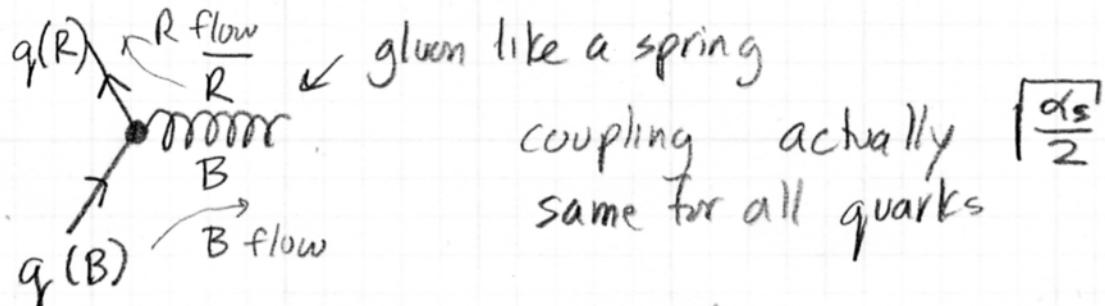


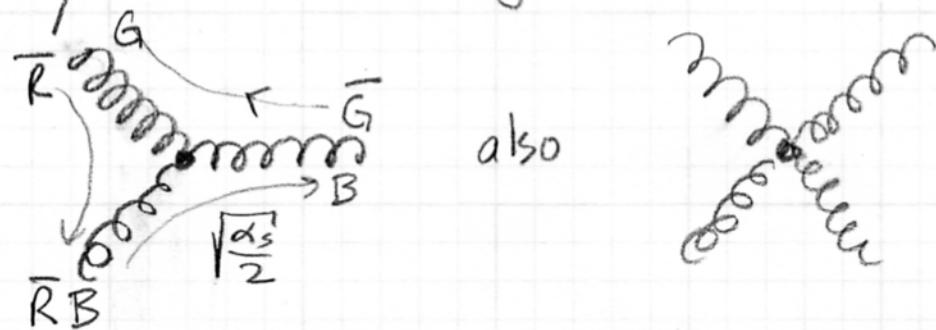
Strong Interaction

- Only quarks & gluons participate
- now each gluon has a color & an anticolor
- Basic vertex:



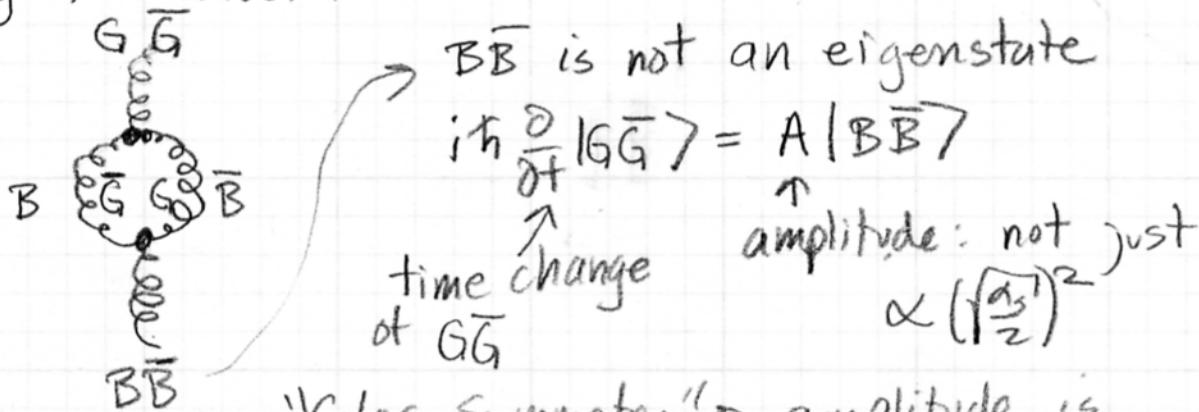
- What about "color neutral" case

Gluons actually couple to themselves, since they have color charge (and anticharge)



α_s generally $\gg 1$, so, the "lowest order" diagram generally greatly underestimates the amplitude. However, the lowest order diagram still gives some guidance.

The first process: where does the "ninth" gluon go? Consider:



"Color Symmetry" \rightarrow amplitude is the same for all color transmutations.

$$\langle R\bar{R} | \psi \rangle = g_1, \quad \langle G\bar{G} | \psi \rangle = g_2, \quad \langle B\bar{B} | \psi \rangle = g_3$$

then $i\hbar \frac{\partial}{\partial t} \begin{pmatrix} g_1 \\ g_2 \\ g_3 \end{pmatrix} = A \begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} g_1 \\ g_2 \\ g_3 \end{pmatrix}$

what are the eigenstates? The Hamiltonian is so simple, with so much symmetry, I'll just "guess"

$$\begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1/r_3 \\ 1/r_3 \\ 1/r_3 \end{pmatrix} = 3 \begin{pmatrix} 1/r_3 \\ 1/r_3 \\ 1/r_3 \end{pmatrix} \quad \text{so } \frac{1}{r_3} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \text{ is eigenvector, eigenvalue } 3A$$

$$\begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1/r_2 \\ -1/r_2 \\ 0 \end{pmatrix} = 0 \begin{pmatrix} 1/r_2 \\ -1/r_2 \\ 0 \end{pmatrix} \quad \text{so } \frac{1}{r_2} \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix} \text{ is eigenvector, eigenvalue } 0$$

$$\begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1/r_2 \\ 1/r_2 \\ -2/r_2 \end{pmatrix} = 0 \begin{pmatrix} 1/r_2 \\ 1/r_2 \\ -2/r_2 \end{pmatrix} \quad \text{so } \frac{1}{r_2} \begin{pmatrix} 1 \\ 1 \\ -2 \end{pmatrix} \text{ is eigenvector, eigenvalue } 0$$

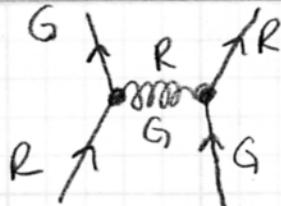
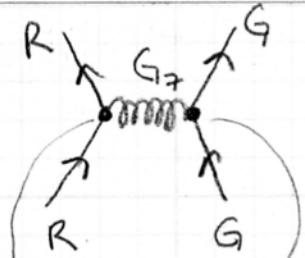
The last two are not unique, because of the degeneracy; those are the "conventional" choices.

For $\frac{1}{r_3} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$, the eigenvalue of energy is $3A$, and $3A \gg$ current energies. So, the "ninth gluon" described by: $|G_9\rangle = \frac{1}{r_2} (|R\bar{R}\rangle + |G\bar{G}\rangle + |B\bar{B}\rangle)$ is so massive it does not participate in low-energy interactions.

The other two, $|G_7\rangle = \frac{1}{r_2} (|R\bar{R}\rangle - |G\bar{G}\rangle)$ and $\frac{1}{r_2} (|R\bar{R}\rangle + |G\bar{G}\rangle - 2|B\bar{B}\rangle) = |G_8\rangle$ have 0 mass, just like $|R\bar{G}\rangle, |R\bar{B}\rangle$, etc

How 2 Quarks Can be Attracted

2 electrons repel via electromagnetic interaction. but 2 quarks of same type can attract by taking advantage of the degree of freedom provided by color. The attraction occurs when the two quarks have different colors.



$$\sqrt{\frac{\alpha_s}{2}} \cdot \frac{1}{r_2} \times \sqrt{\frac{\alpha_s}{2}} \times -\frac{1}{r_2} = -\frac{\alpha_s}{4}$$

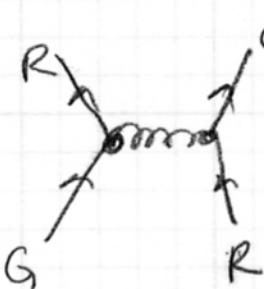
$$\sqrt{\frac{\alpha_s}{2}} \cdot \frac{1}{r_6} \times \sqrt{\frac{\alpha_s}{2}} \cdot \frac{1}{r_6}$$

$$\frac{\alpha_s}{2}$$

whence,
 $q(R)q(G)$ not
 an
 eigenstate!!

$$+\frac{\alpha_s}{12}$$

$$-\frac{1}{4} + \frac{1}{r_2} = -\frac{1}{6}$$



$$a_1 = \langle RG | \psi \rangle \quad a_2 = \langle GR | \psi \rangle$$

$$\text{and it } \frac{\partial}{\partial t} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \alpha_s \begin{pmatrix} -1/6 & 1/2 \\ 1/2 & -1/6 \end{pmatrix} \in \text{very symmetric.}$$

c eigenvectors by guess: $\begin{pmatrix} -1/6 & 1/2 \\ 1/2 & -1/6 \end{pmatrix} \begin{pmatrix} 1/r_2 \\ 1/r_2 \end{pmatrix} = \frac{1}{6} \begin{pmatrix} 1/r_2 \\ 1/r_2 \end{pmatrix}$ eigenvalue $\frac{\alpha_s}{6}$

$\begin{pmatrix} -1/6 & 1/2 \\ 1/2 & -1/6 \end{pmatrix} \begin{pmatrix} -1/r_2 \\ -1/r_2 \end{pmatrix} = -\frac{2}{3} \begin{pmatrix} 1/r_2 \\ 1/r_2 \end{pmatrix}$ eigenvalue $-\frac{2}{3} \alpha_s$

The punchline: $\frac{1}{r_2} (\langle q(R)q(B) \rangle - \langle q(B)q(R) \rangle)$

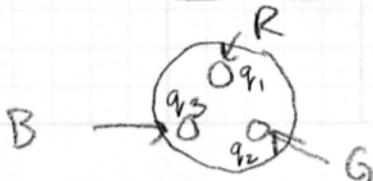
"asymmetric" has negative (attractive) energy,
 magnitude $-\frac{2}{3} \alpha_s$

Proton/Neutron/Baryon: $c_1 = R \quad c_2 = G \quad c_3 = B$

$$| \text{Baryon} \rangle \propto \sum \epsilon_{ijk} | q_1(c_i) q_2(c_j) q_3(c_k) \rangle$$

$\epsilon_{ijk} = 0$ if any 2 indices same

$$\epsilon_{123} = \epsilon_{312} = \epsilon_{231} = 1; \quad \epsilon_{213} = \epsilon_{321} = \epsilon_{132} = -1$$

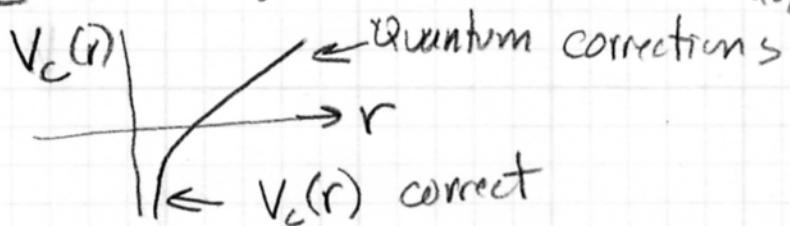


So, if you exchange any pair of colors (or, separately, quarks) is swapped, the state comes back to itself, multiplied by -1 ... "totally asymmetric!"

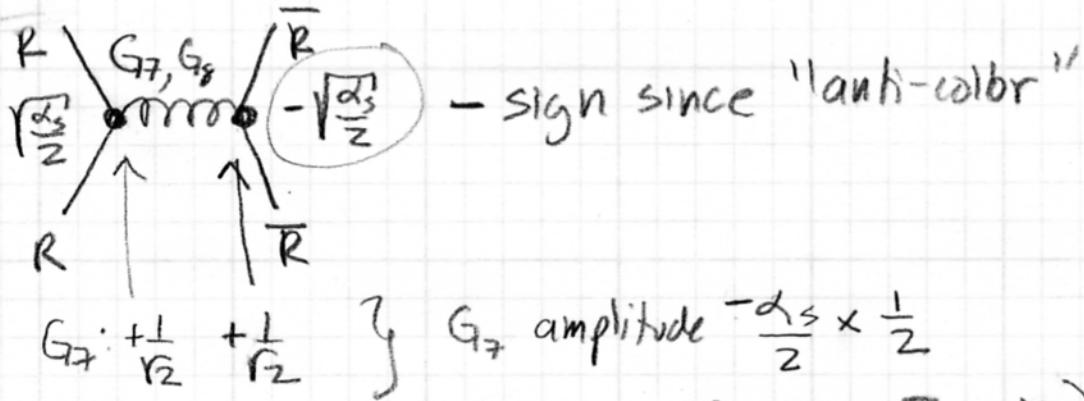
\Rightarrow "TOTALLY ATTRACTIVE" (you are held together)

note: • spatial dependence of strong potential has been ignored

- at short distance, $V_c(r) = -\frac{2}{3} \frac{\alpha_s}{r} q\bar{q}$ "asymmetric"



- How about $q\bar{q}$? Most interesting when these are colorless

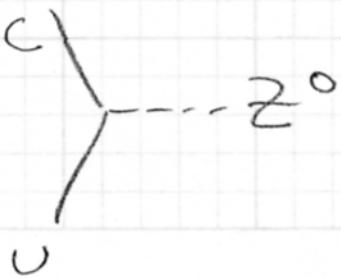


you do G_8 , + others ($R\bar{R} \rightarrow G\bar{G}$ etc)
ignore annihilation diagrams, only
do "+-channel"

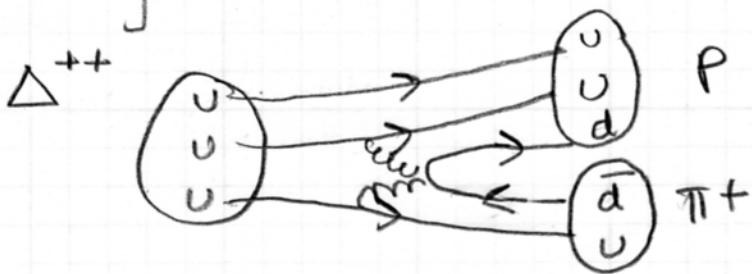
Other comments about decays, Feynman diagrams:

1) Z^0 like a photon; never

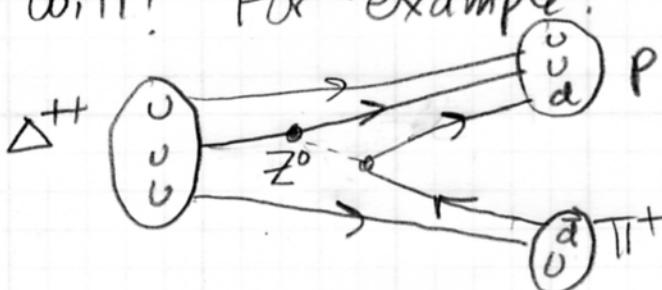
BIG MISTERY.



2) Strong interaction can drive decays.



- * no quark changes flavor in strong decays
- * quark-antiquark pair often "pops" from glue - usually $d\bar{d}$ or $u\bar{u}$; rarely $s\bar{s}$; $c\bar{c}, b\bar{b}, t\bar{t}$ too heavy, suppressed
- * if a decay can go strong, it will! For example:

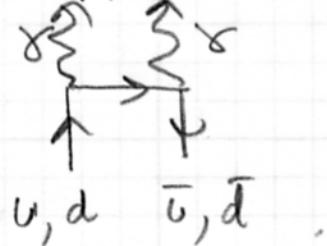


can go,
but
Branching
ratio negligibl.

3) Symmetry between $u + d$ can complicate strong decays.. "isospin"

4) Sometimes strong can't go... Then, electromagnetic tries

$$\pi^0 \rightarrow \gamma\gamma$$



always γ 's
or e^+e^- in
final state.

5) when strong em can't go,
Weak can (changing flavor of quarks).