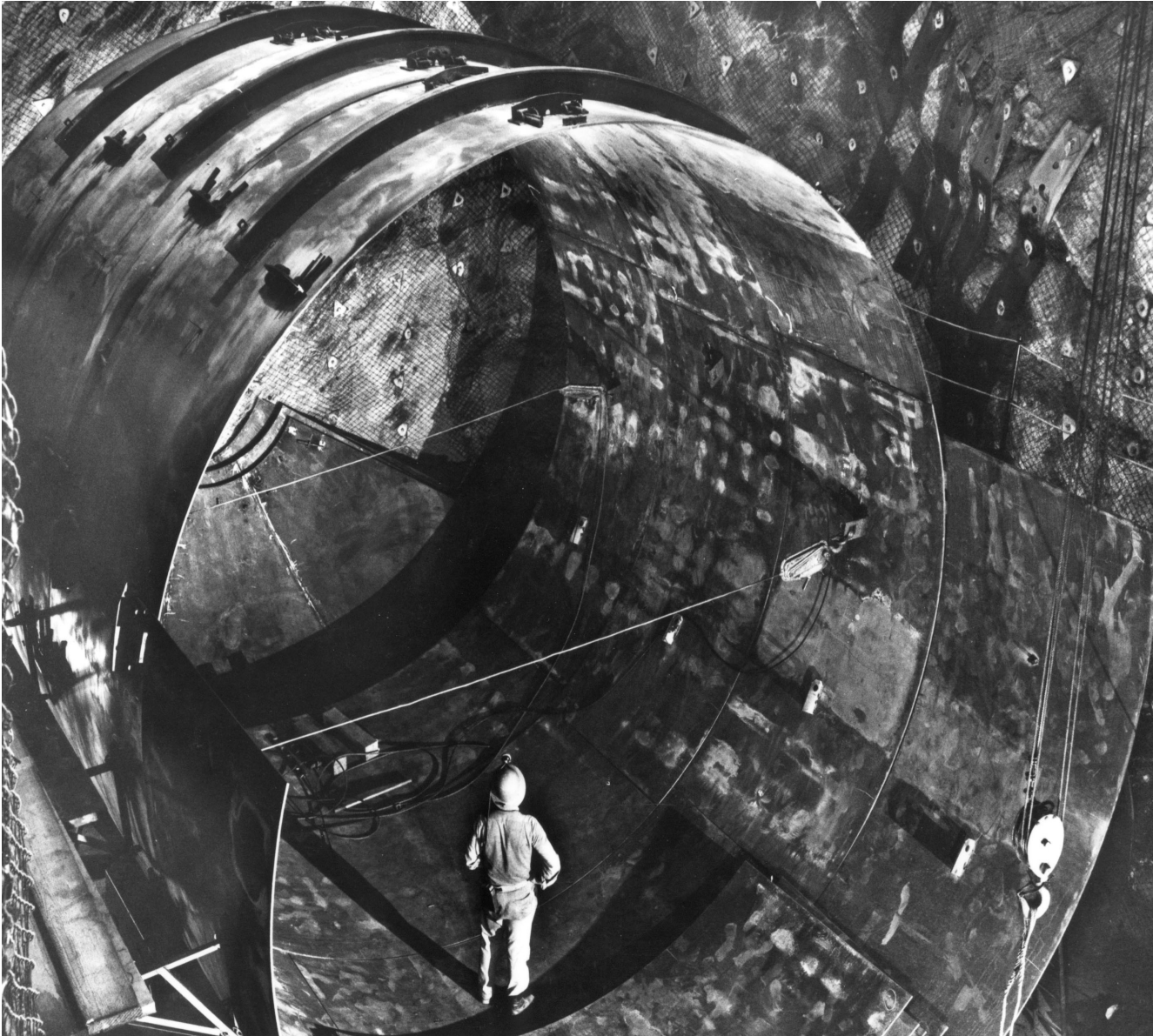


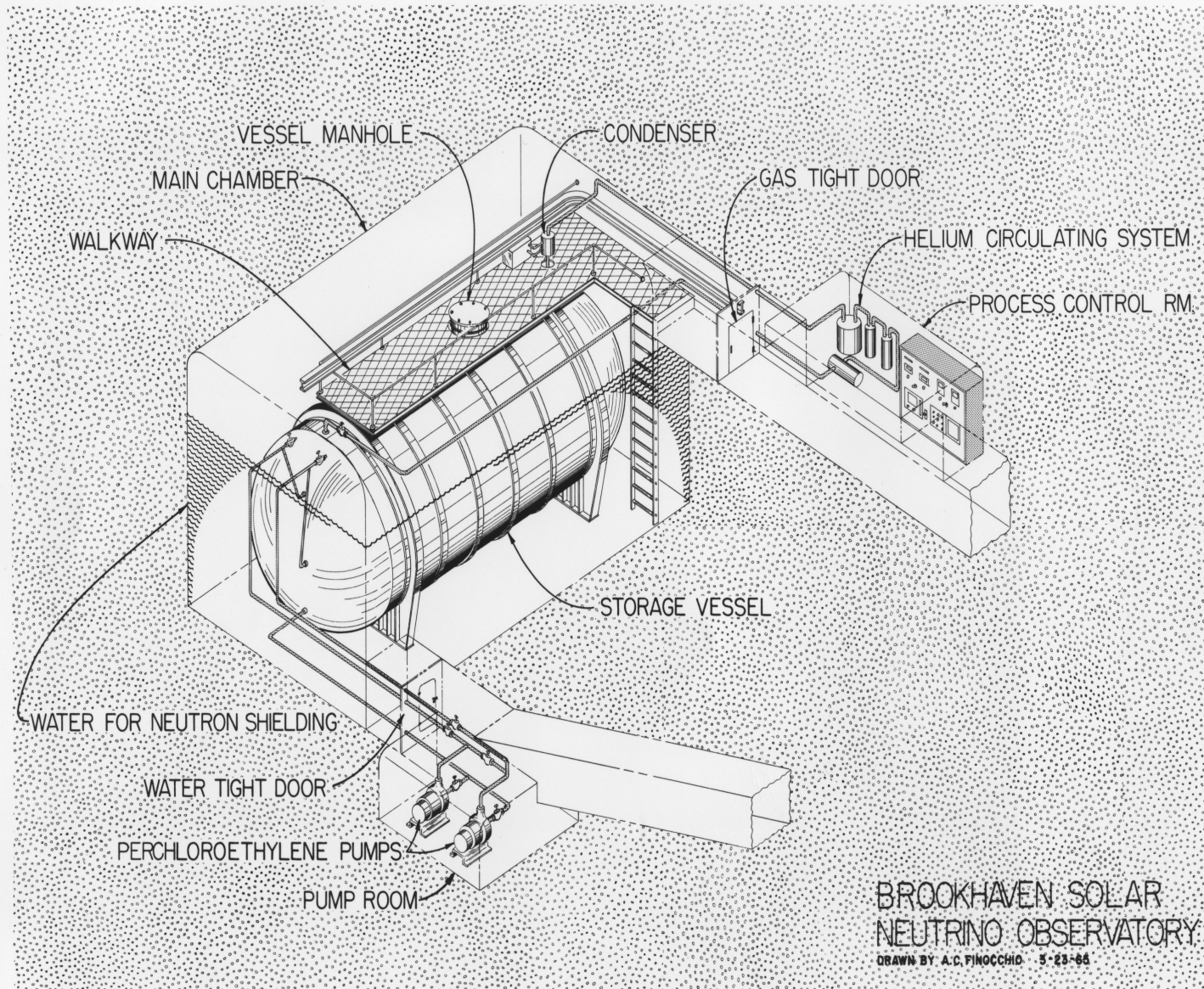
# Physics 125

5/10/2010

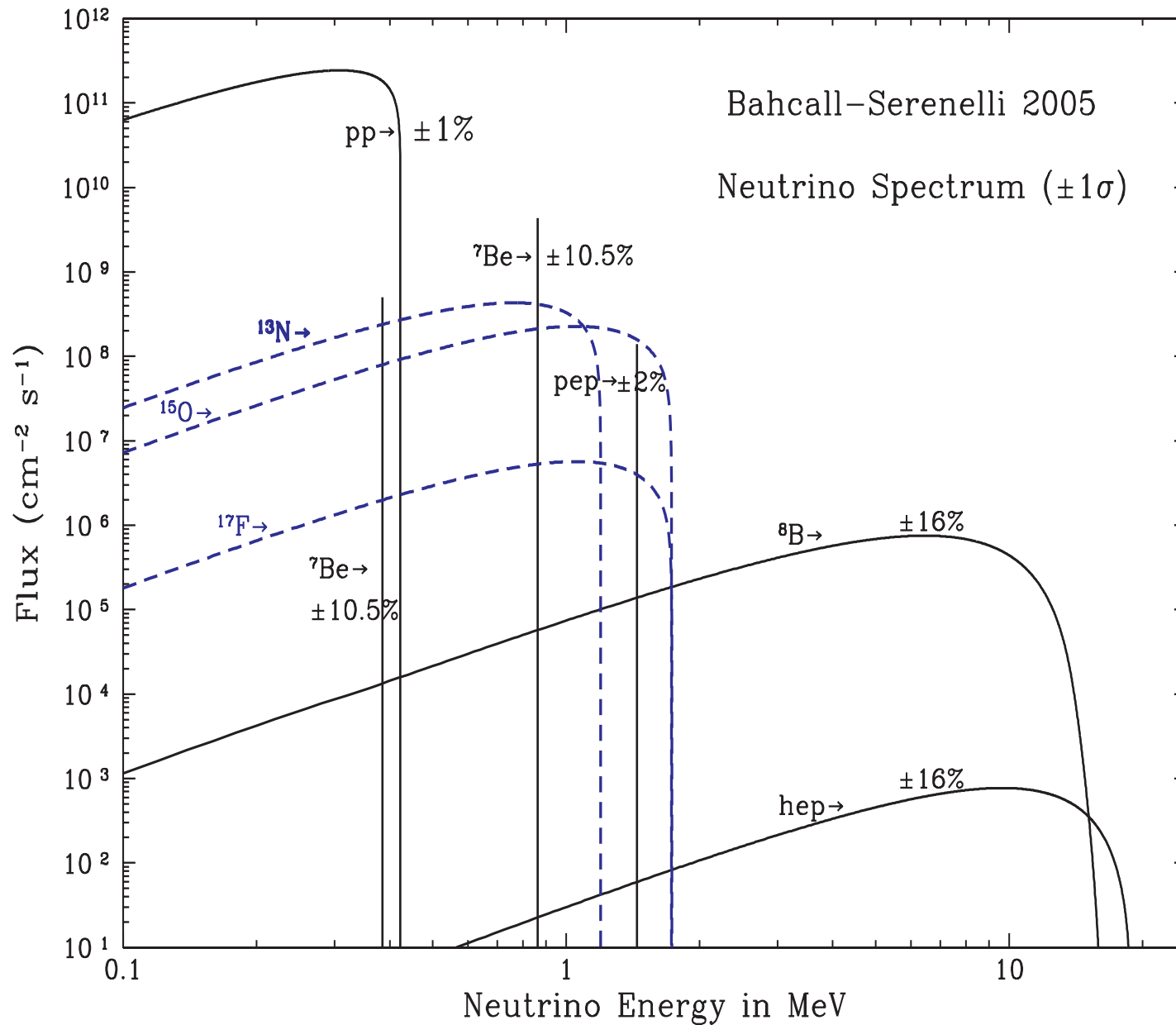
# Davis Tank – Near Deadwood



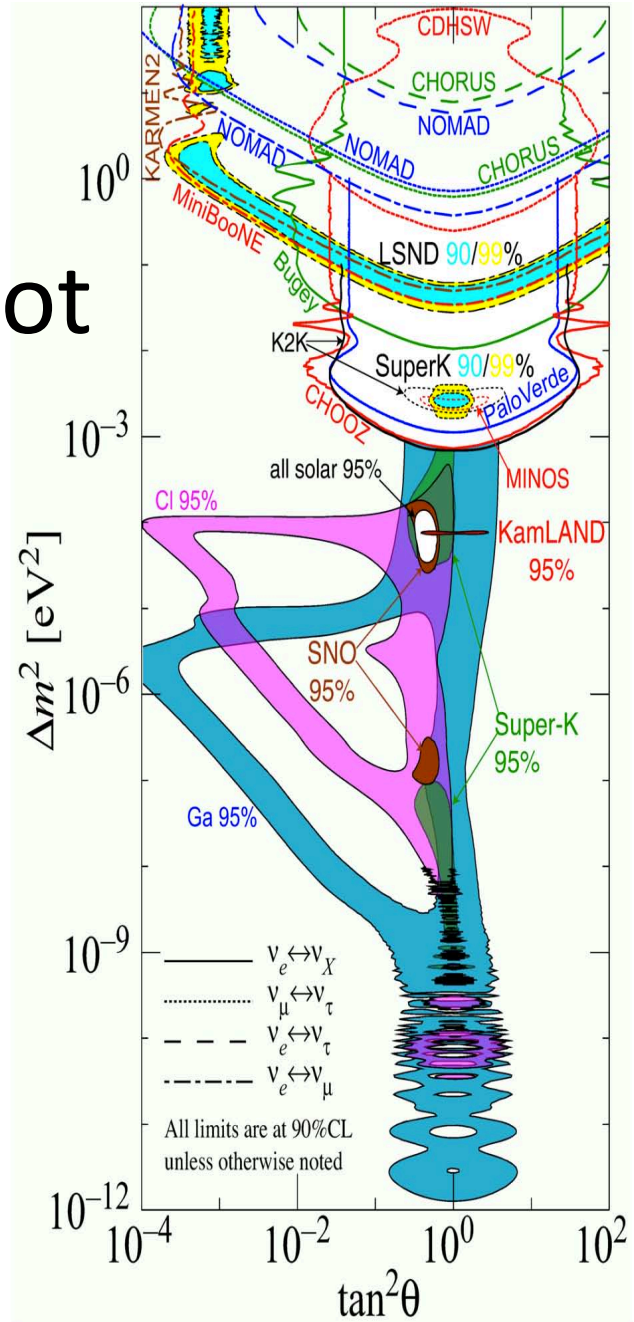
# Davis Experiment



# Solar Neutrino Flux

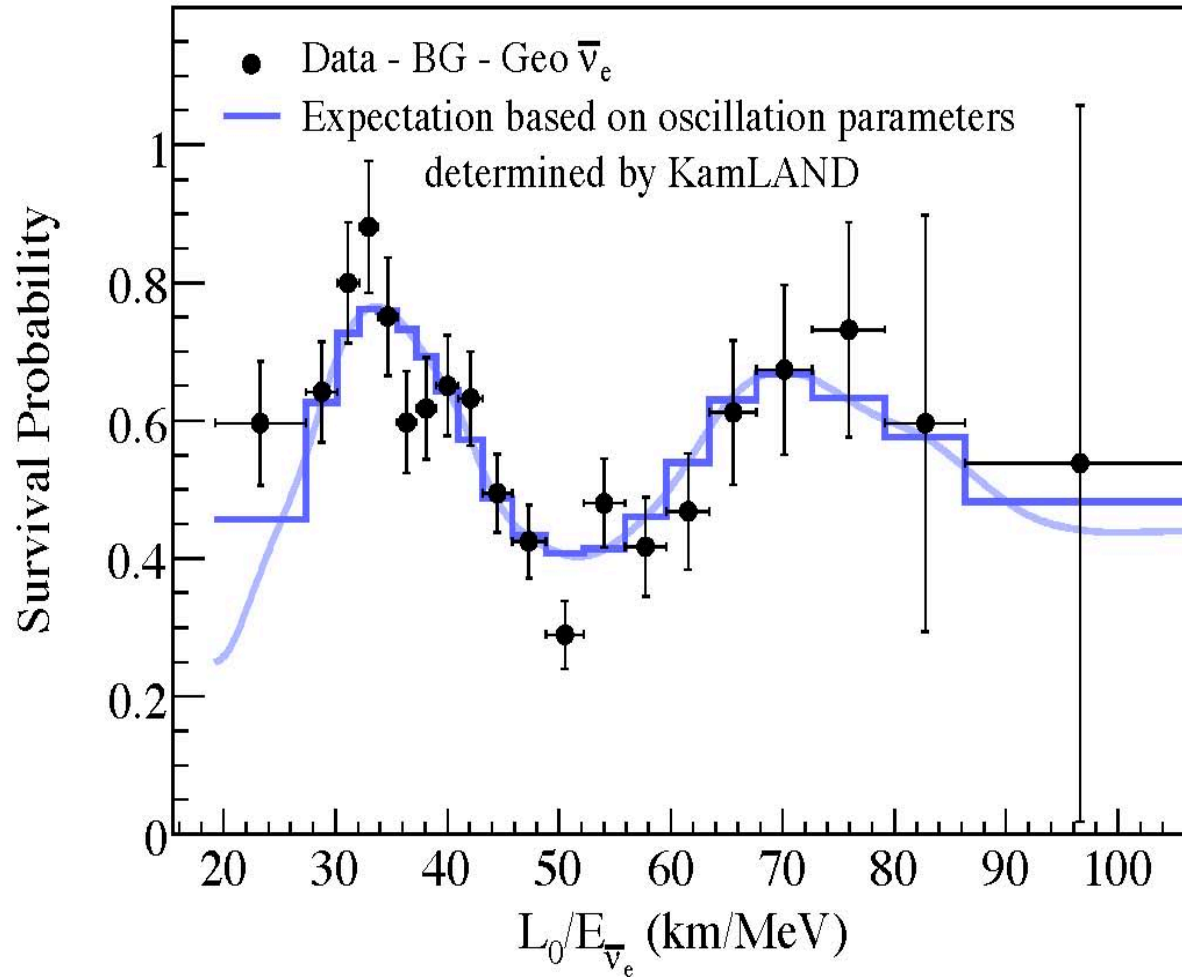


# The Search Plot



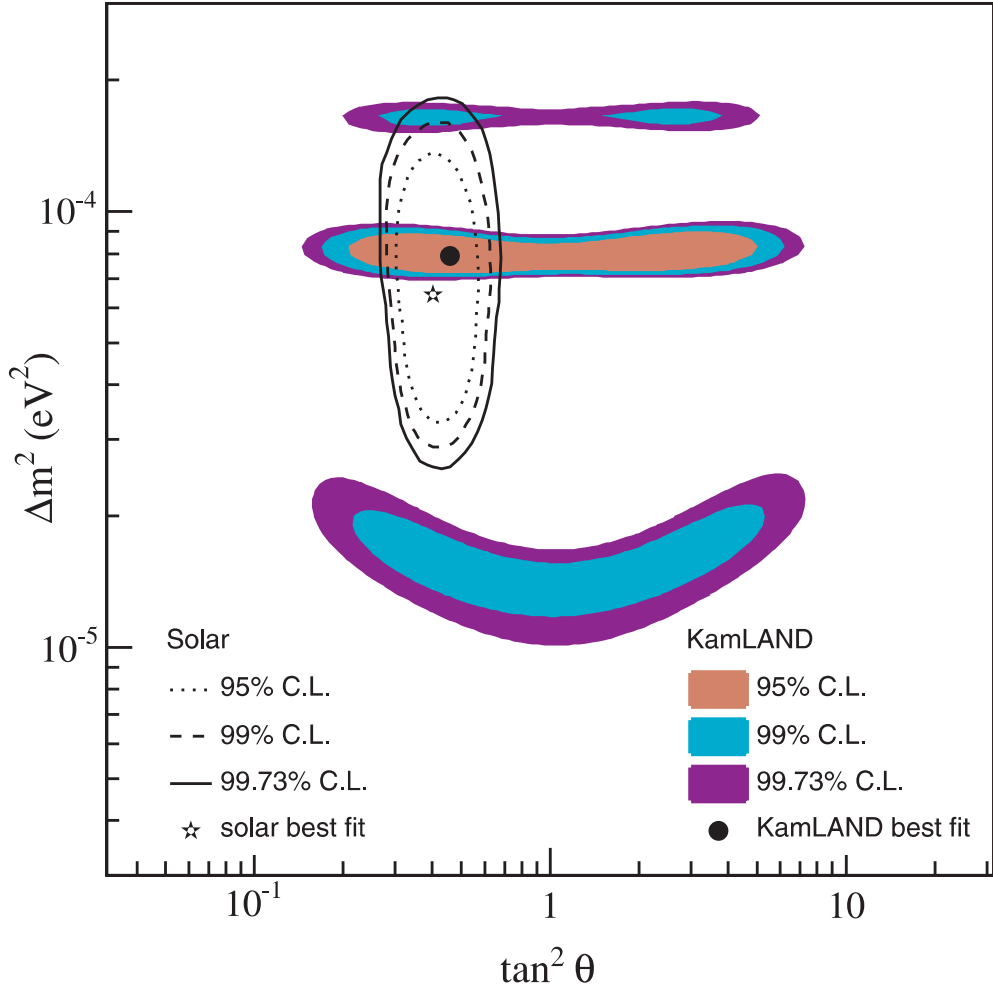
**Figure 13.4:** The regions of squared-mass splitting and mixing angle favored or excluded by various experiments. This figure was contributed by H. Murayama (University of California, Berkeley). References to the data used in the figure can be found at <http://hitoshi.berkeley.edu/neutrino/>.

# Earth-based reactor experiment

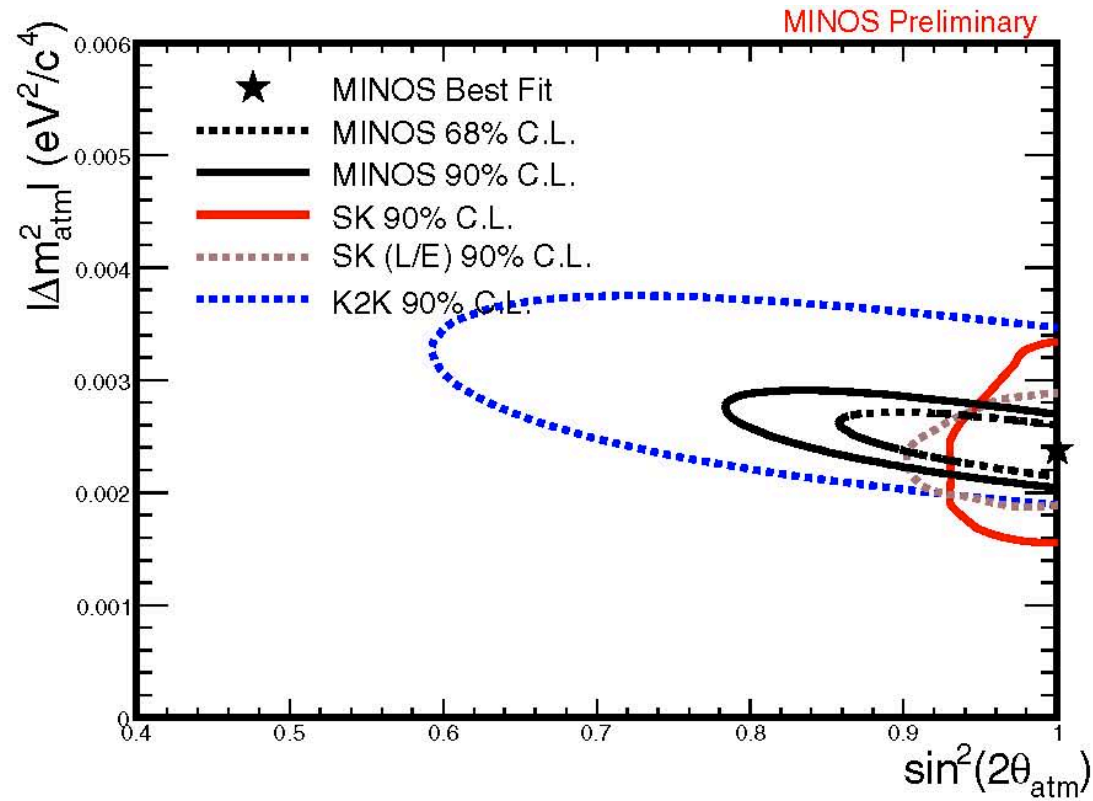


**Figure 13.2:** Ratio of the background- and geo-neutrino subtracted  $\bar{\nu}_e$  spectrum to the no-oscillation expectation as a function of  $L_0/E$  [28]. See text for explanation.

# Solar/KAMLAND



# Atmospheric...



**Figure 13.1:** The region of the atmospheric oscillation parameters  $\Delta m_{\text{atm}}^2$  and  $\sin^2 2\theta_{\text{atm}}$  allowed by the SK, K2K, and MINOS data. The results of two different analyses of the SK (“Super K”) data are shown [21].



# Electron Neutrino Appearance (RPP)

With a conventional beam, one would seek  $CP$  violation, and try to determine whether the mass spectrum is normal or inverted, by studying the oscillations  $\nu_\mu \rightarrow \nu_e$  and  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ . The appearance probability for  $\nu_e$  in a beam that is initially  $\nu_\mu$  can be written for  $\sin^2 2\theta_{13} < 0.2$  [53]

$$P(\nu_\mu \rightarrow \nu_e) \cong \sin^2 2\theta_{13} T_1 - \alpha \sin 2\theta_{13} T_2 + \alpha \sin 2\theta_{13} T_3 + \alpha^2 T_4 \quad . \quad (13.39)$$

Here,  $\alpha \equiv \Delta m_{21}^2 / \Delta m_{31}^2$  is the small ( $\sim 1/30$ ) ratio between the solar and atmospheric squared-mass splittings, and

$$T_1 = \sin^2 \theta_{23} \frac{\sin^2[(1-x)\Delta]}{(1-x)^2} \quad , \quad (13.40)$$

$$T_2 = \sin \delta \sin 2\theta_{12} \sin 2\theta_{23} \sin \Delta \frac{\sin(x\Delta)}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \quad , \quad (13.41)$$

$$T_3 = \cos \delta \sin 2\theta_{12} \sin 2\theta_{23} \cos \Delta \frac{\sin(x\Delta)}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \quad , \quad (13.42)$$

and

$$T_4 = \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(x\Delta)}{x^2} \quad . \quad (13.43)$$

In these expressions,  $\Delta \equiv \Delta m_{31}^2 L / 4E$  is the kinematical phase of the oscillation. The quantity  $x \equiv 2\sqrt{2}G_F N_e E / \Delta m_{31}^2$ , with  $G_F$  the Fermi coupling constant and  $N_e$  the electron number density, is a measure of the importance of the matter effect resulting from coherent forward-scattering of electron neutrinos from ambient electrons as the neutrinos travel through the earth from the source to the detector [cf. Sec. I]. In the appearance probability  $P(\nu_\mu \rightarrow \nu_e)$ , the  $T_1$  term represents the oscillation due to the atmospheric-mass-splitting scale, the  $T_4$  term represents the oscillation due to the solar-mass-splitting scale, and the  $T_2$  and  $T_3$  terms are the  $CP$ -violating and  $CP$ -conserving interference terms, respectively.

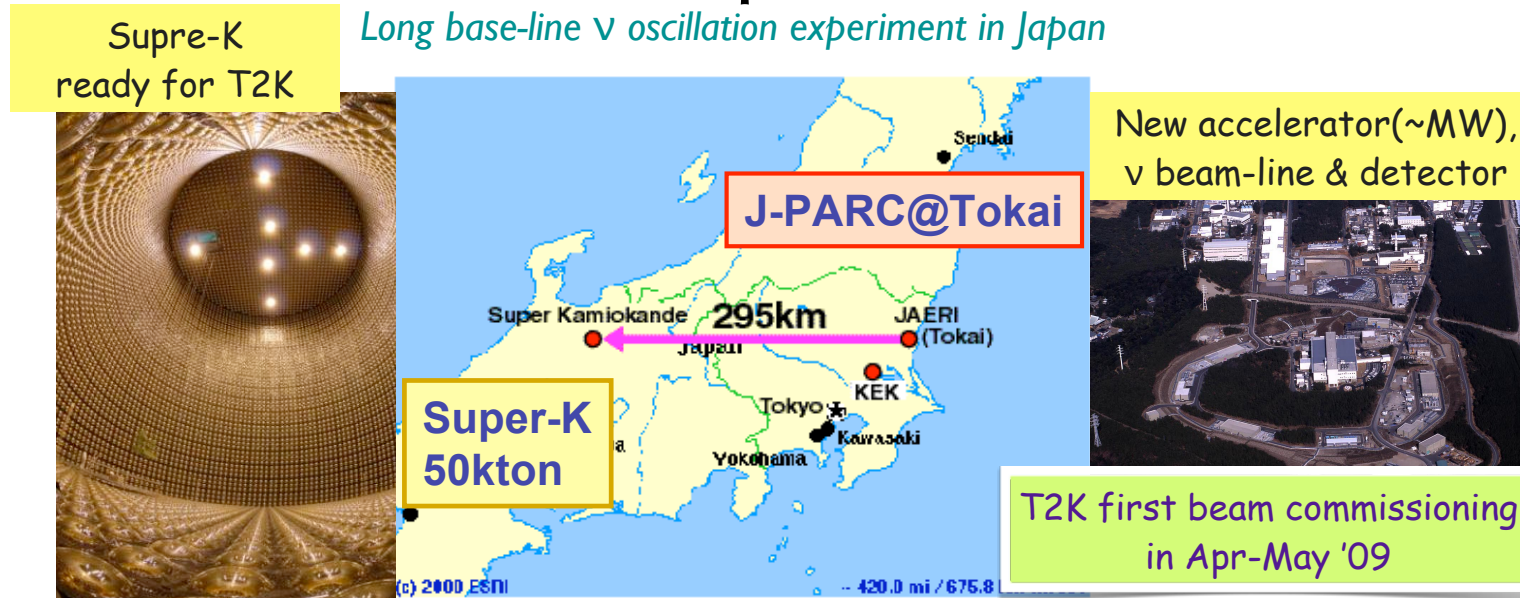
The probability for the corresponding antineutrino oscillation,  $P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)$ , is the same as the probability  $P(\nu_\mu \rightarrow \nu_e)$  given by Eqs. (13.39)–(13.43), but with the signs in front

# Electron Neutrino Appearance (2) (RPP)

of both  $x$  and  $\sin \delta$  reversed: both the matter effect and  $CP$  violation lead to a difference between the  $\nu_\mu \rightarrow \nu_e$  and  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillation probabilities. In view of the dependence of  $x$  on  $\Delta m_{31}^2$ , and in particular on the sign of  $\Delta m_{31}^2$ , the matter effect can reveal whether the neutrino mass spectrum is normal or inverted. However, to determine the nature of the spectrum, and to establish the presence of  $CP$  violation, it obviously will be necessary to disentangle the matter effect from  $CP$  violation in the neutrino-antineutrino oscillation probability difference that is actually observed. To this end, complementary measurements will be extremely important. These can take advantage of the differing dependences on the matter effect and on  $CP$  violation in  $P(\nu_\mu \rightarrow \nu_e)$ .

# Tokai to Kamiokande - Sakashita

## T2K Experiment



### T2K features to enhance the sensitivity

- ▶ **Super-K(SK) as main neutrino detector** :  
22.5kton(fiducial) water cherenkov detector & good PID ( $e/\mu$ ) performance
- ▶ **Off-axis beam** (intense & narrow-band low energy neutrino beam) → next slide
- ▶ **Neutrino energy reconstruction** :  
CCQE interactions dominate at T2K beam energy

# Tokai to Kamiokande - Sakashita

Next step of  $\nu$  oscillation experiment

Neutrino Oscillation Experiment

- discover a finite  $\theta_{13}$ 
  - determine  $|U_{e3}|$ 
    - important role for future neutrino experiments
    - CPV in lepton sector
      - hint on Baryon# asymmetry of Universe
    - mass hierarchy
- precise measurement
  - Is  $\theta_{23}$  maximal ?

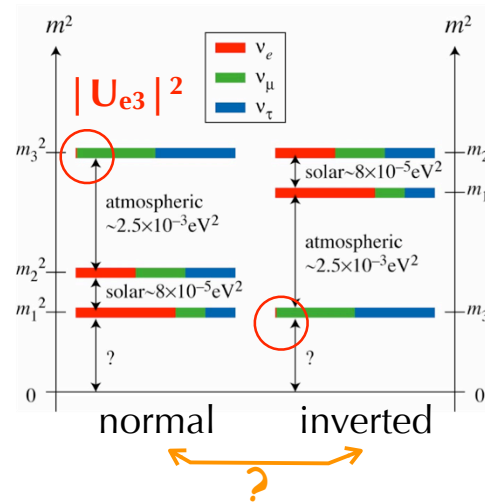
$$U_{MNSP} \sim \begin{pmatrix} 0.8 & 0.5 & \boxed{?} \\ 0.4 & 0.6 & 0.7 \\ 0.4 & 0.6 & 0.7 \end{pmatrix} \quad U_{e3} = s_{13}e^{-i\delta}$$

$0\nu\beta\beta$   
decay exp.

- Dirac or Majorana

Tritium  $\beta$   
decay exp.

- absolute mass scale



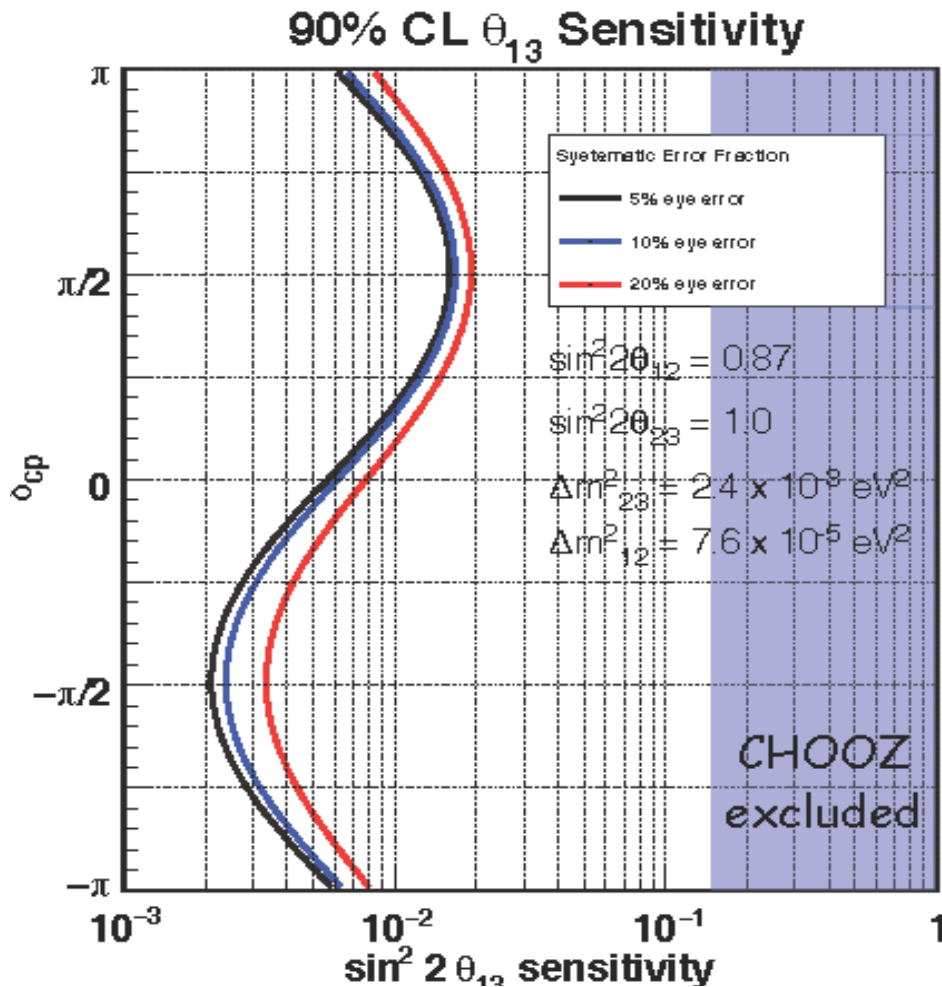
# Nova – starts about 2013



# The 'Money Plots'

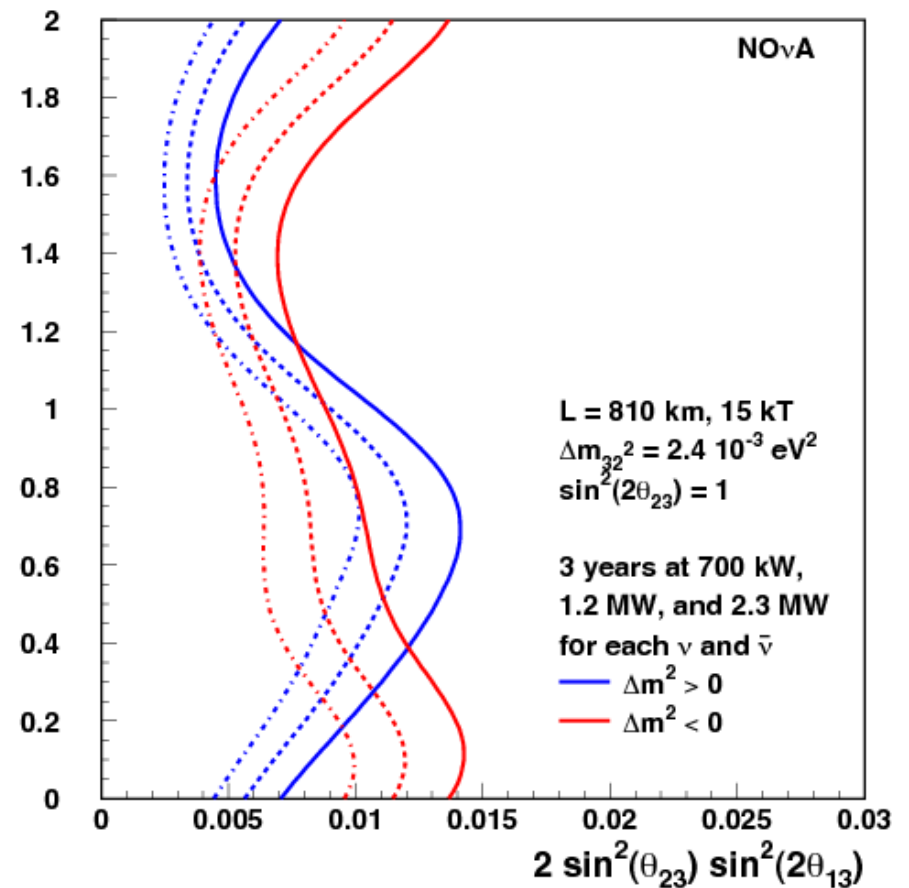
T2K

> x10 improvement from CHOOZ limit

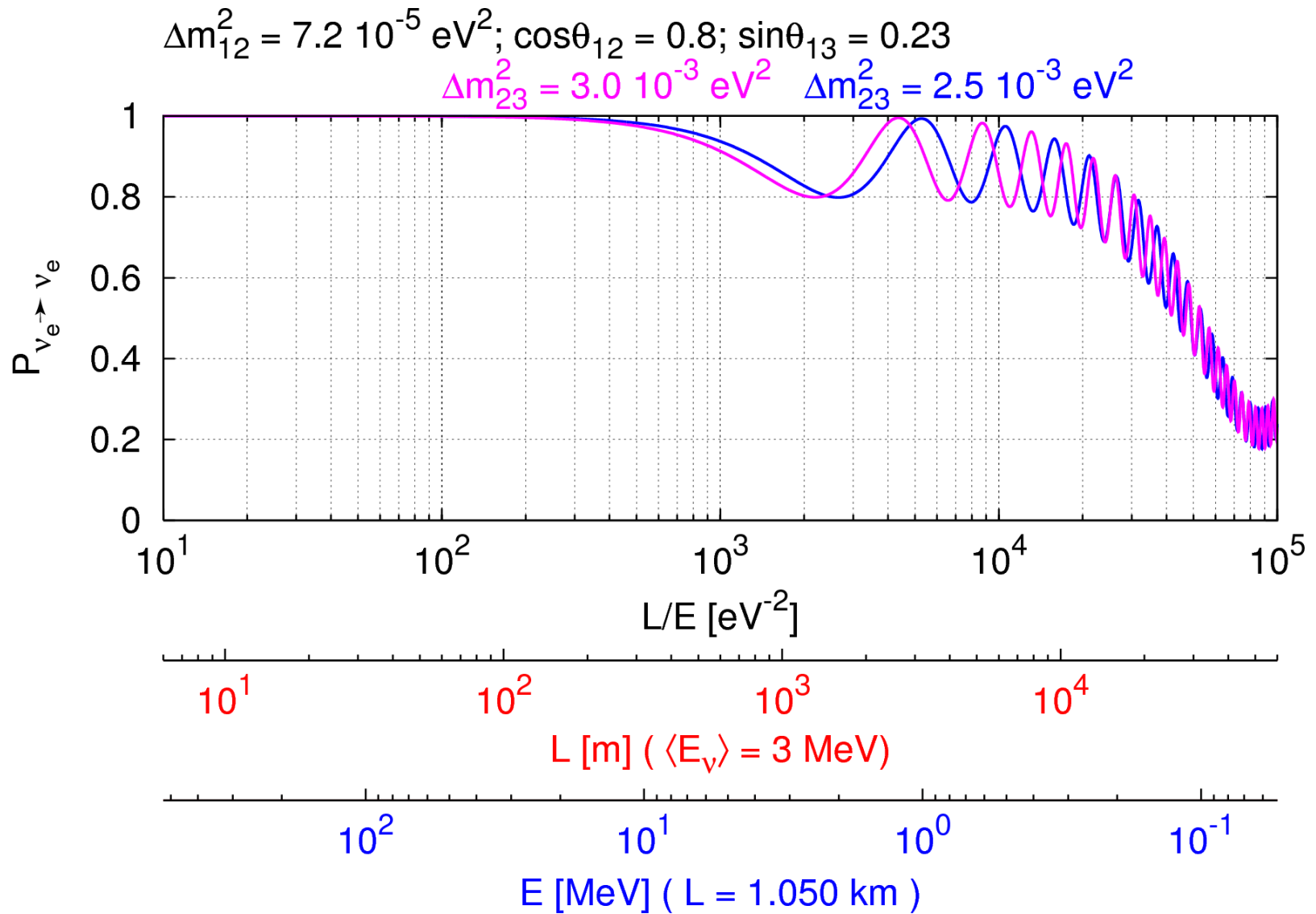


Nova

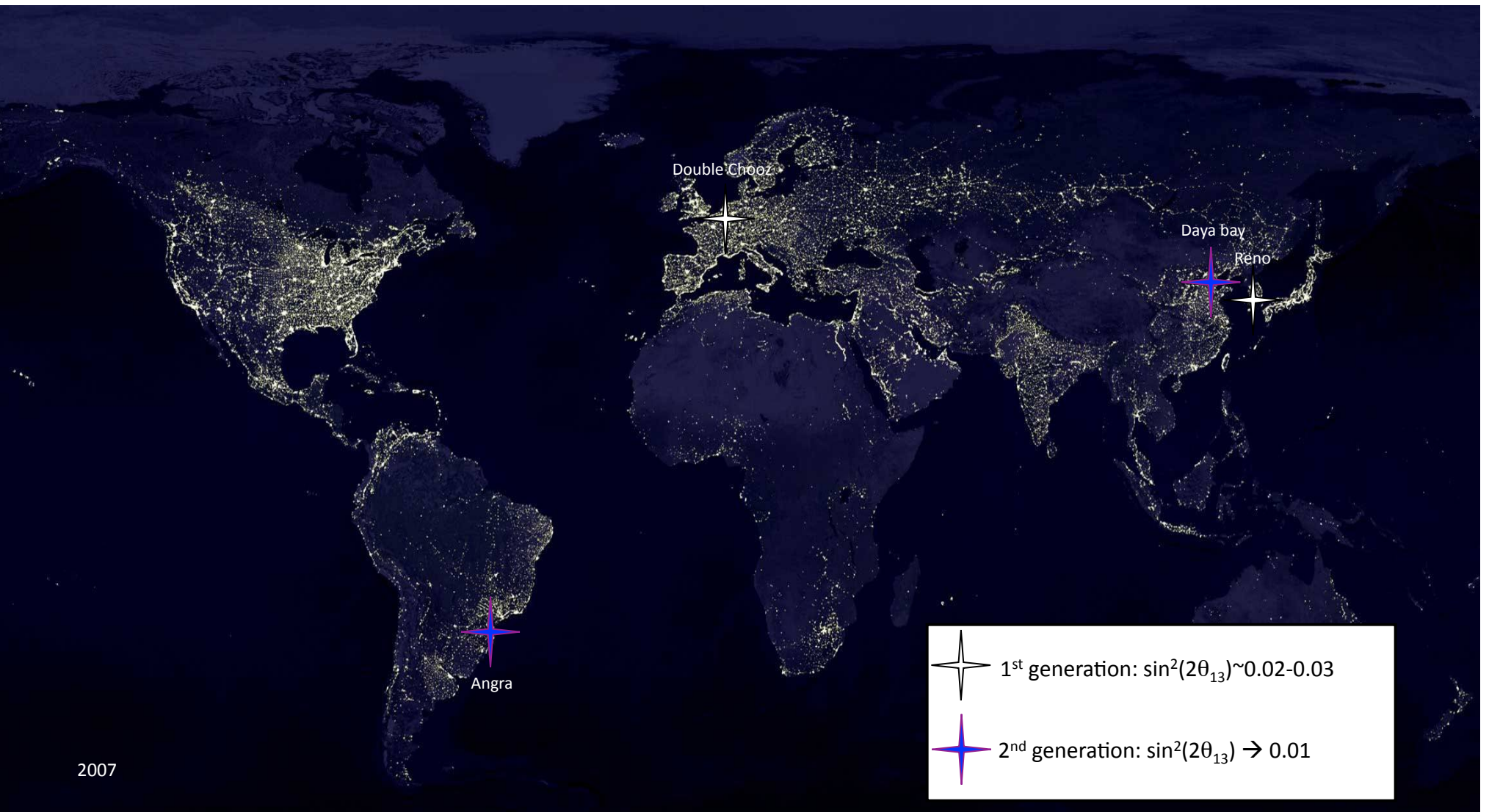
90% CL Sensitivity to  $\sin^2(2\theta_{13}) \neq 0$



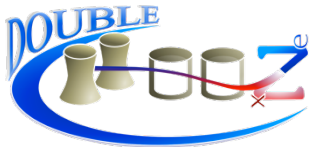
# Electron antineutrino disappearance



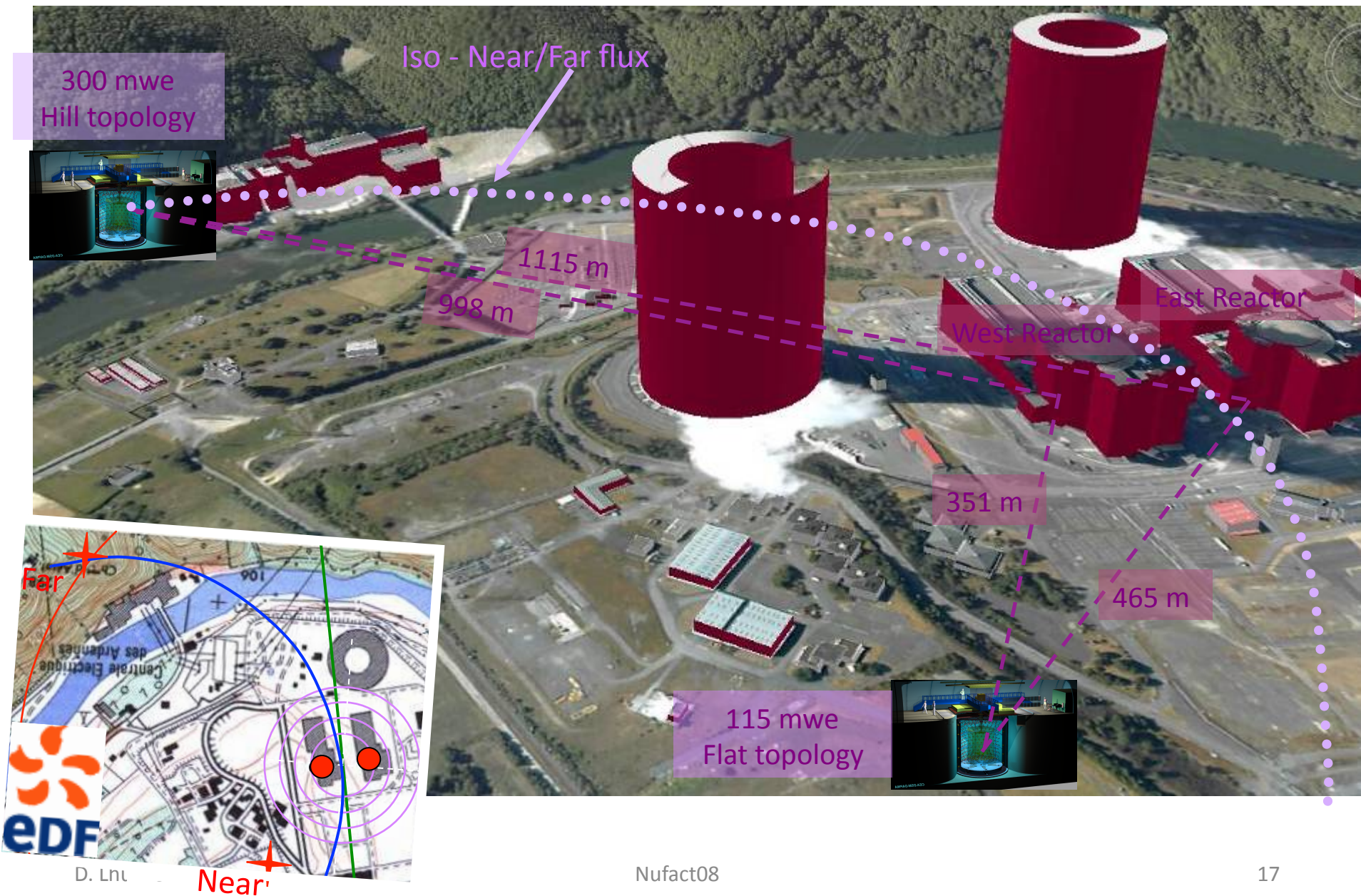
# Looking for Sites







# Site in French Ardennes



# Daya Bay...



# Perpectives @ Reactors

P.Huber, M.Lindner, T.Schwetz and W.Winter

Giant detectors?

Nucl.Phys.B665:487-519,2003

