

The idea is to factor out the expected behavior in $U_{\tilde{\epsilon}l}(\tilde{\rho})$. It's easiest to do this in 2 steps:

$$\#1 \quad U_{\tilde{\epsilon}l}(\tilde{\rho}) = e^{-\alpha\tilde{\rho}} v_{\tilde{\epsilon}l}(\tilde{\rho}) \quad \left\{ \begin{array}{l} \text{as } \tilde{\rho} \rightarrow \infty \\ v_{\tilde{\epsilon}l} \text{ must} \\ \text{not spoil } e^{-\alpha\tilde{\rho}}! \end{array} \right.$$

$$U'_{\tilde{\epsilon}l}(\tilde{\rho}) = -\alpha e^{-\alpha\tilde{\rho}} v_{\tilde{\epsilon}l}(\tilde{\rho}) + e^{-\alpha\tilde{\rho}} v'_{\tilde{\epsilon}l}(\tilde{\rho})$$

$$U''_{\tilde{\epsilon}l}(\tilde{\rho}) = \alpha^2 e^{-\alpha\tilde{\rho}} v_{\tilde{\epsilon}l}(\tilde{\rho}) - 2\alpha e^{-\alpha\tilde{\rho}} v'_{\tilde{\epsilon}l}(\tilde{\rho}) + e^{-\alpha\tilde{\rho}} v''_{\tilde{\epsilon}l}(\tilde{\rho})$$

plugging back in:

$$\alpha^2 e^{-\alpha\tilde{\rho}} v_{\tilde{\epsilon}l}(\tilde{\rho}) + 2\alpha e^{-\alpha\tilde{\rho}} v'_{\tilde{\epsilon}l}(\tilde{\rho}) - e^{-\alpha\tilde{\rho}} v''_{\tilde{\epsilon}l}(\tilde{\rho})$$

$$+ \frac{l(l+1)}{\tilde{\rho}^2} e^{-\alpha\tilde{\rho}} v_{\tilde{\epsilon}l}(\tilde{\rho}) - \frac{2}{\tilde{\rho}} e^{-\alpha\tilde{\rho}} v'_{\tilde{\epsilon}l}(\tilde{\rho}) = \tilde{\epsilon} e^{-\alpha\tilde{\rho}} v_{\tilde{\epsilon}l}(\tilde{\rho})$$

these cancel if we choose $\tilde{\epsilon} = -\alpha^2$
since this is a bound state, $\tilde{\epsilon} < 0$

$$\text{and } \boxed{\alpha = \sqrt{-\tilde{\epsilon}}} *$$

cancelling $e^{-\alpha\tilde{\rho}}$

$$\# \quad v''_{\tilde{\epsilon}l}(\tilde{\rho}) - 2\alpha v'_{\tilde{\epsilon}l}(\tilde{\rho}) + \left[\frac{2}{\tilde{\rho}} - \frac{l(l+1)}{\tilde{\rho}^2} \right] v_{\tilde{\epsilon}l}(\tilde{\rho}) = 0$$

in general, the next step is to take

$$v_{\tilde{\epsilon}l}(\tilde{\rho}) = \tilde{\rho}^{l+1} w_{\tilde{\epsilon}l}(\tilde{\rho})$$

$$\text{and } w_{\tilde{\epsilon}l}(\tilde{\rho}) = \sum_{k=0}^{\infty} C_k \tilde{\rho}^k \quad \leftarrow \begin{array}{l} k=0 \text{ start} \\ \text{because of } \tilde{\rho}^{l+1} \end{array}$$

but look out: C_k must not spoil $e^{-\alpha\tilde{\rho}}$ suppression

$$V'_{\hat{\epsilon}_l}(\tilde{\rho}) = (l+1)\tilde{\rho}^l W_{\hat{\epsilon}_l}(\tilde{\rho}) + \tilde{\rho}^{l+1} W'_{\hat{\epsilon}_l}(\tilde{\rho})$$

$$V''_{\hat{\epsilon}_l}(\tilde{\rho}) = l(l+1)\tilde{\rho}^{l-1} W_{\hat{\epsilon}_l}(\tilde{\rho}) + 2(l+1)\tilde{\rho}^l W'_{\hat{\epsilon}_l}(\tilde{\rho}) + \tilde{\rho}^{l+1} W''_{\hat{\epsilon}_l}(\tilde{\rho})$$

plugging back into the ~~*~~ equation,

$$l(l+1)\tilde{\rho}^{l-1} W_{\hat{\epsilon}_l}(\tilde{\rho}) + 2(l+1)\tilde{\rho}^l W'_{\hat{\epsilon}_l}(\tilde{\rho}) + \tilde{\rho}^{l+1} W''_{\hat{\epsilon}_l}(\tilde{\rho}) - 2\alpha \left\{ (l+1)\tilde{\rho}^l W_{\hat{\epsilon}_l}(\tilde{\rho}) + \tilde{\rho}^{l+1} W'_{\hat{\epsilon}_l}(\tilde{\rho}) \right\} + 2\tilde{\rho}^l W_{\hat{\epsilon}_l}(\tilde{\rho}) - l(l+1)\tilde{\rho}^{l-1} W_{\hat{\epsilon}_l}(\tilde{\rho}) = 0 \quad \del{*}$$

now $W_{\hat{\epsilon}_l}(\tilde{\rho}) = \sum_{k=0}^{\infty} C_k \tilde{\rho}^k$

$$W'_{\hat{\epsilon}_l}(\tilde{\rho}) = \sum_{k=0}^{\infty} C_k k \tilde{\rho}^{k-1} = \sum_{k=0}^{\infty} C_{k+1} (k+1) \tilde{\rho}^k$$

$$W''_{\hat{\epsilon}_l}(\tilde{\rho}) = \sum_{k=0}^{\infty} C_{k+1} k(k+1) \tilde{\rho}^{k-1}$$

plug in ~~*~~

$$2(l+1)\tilde{\rho}^l \left[\sum_{k=0}^{\infty} C_{k+1} (k+1) \tilde{\rho}^k \right] + \tilde{\rho}^{l+1} \left[\sum_{k=0}^{\infty} C_{k+1} k(k+1) \tilde{\rho}^{k-1} \right]$$

$$- 2\alpha \left\{ (l+1)\tilde{\rho}^l \sum_{k=0}^{\infty} C_k \tilde{\rho}^k + \tilde{\rho}^{l+1} \sum_{k=0}^{\infty} C_k k \tilde{\rho}^{k-1} \right\}$$

$$+ 2\tilde{\rho}^l \sum_{k=0}^{\infty} C_k \tilde{\rho}^k = 0$$

note that as written, every term is of order $\tilde{\rho}^{l+k}$, so we can collect the coefficients.

$$\left(\underbrace{2(l+1)(k+1) + k(k+1)}_{=(k+1)(2(l+1)+k)} \right) C_{k+1} - 2(\alpha[l+1+k] - 1) C_k = 0$$

$$\text{so, } \boxed{\frac{C_{k+1}}{C_k} = 2 \frac{\alpha[l+1+k] - 1}{(k+1)(k+2(l+1))}}$$

- 1) Can an infinite series work?
In that case,

$$\lim_{k \rightarrow \infty} \frac{C_{k+1}}{C_k} = \frac{2\alpha}{k} \text{ like } e^{+2\alpha\tilde{p}}$$

$$= 1 + 2\alpha\tilde{p} + \frac{1}{2}(2\alpha\tilde{p})^2 + \frac{1}{3!}(2\alpha\tilde{p})^3 + \dots$$

$$\text{thus } U_{\tilde{\epsilon}l} \propto \tilde{p}^{l+1} \underbrace{e^{-\alpha\tilde{p}} e^{+2\alpha\tilde{p}}}_{\text{diverges - recall numerical solution}}$$

- 2) series will terminate if for some k :

$$\frac{C_{k+1}}{C_k} = 0 = \alpha[l+1+k] - 1$$

$$\alpha = \sqrt{-\tilde{\epsilon}} = \frac{1}{k+l+1}$$

$$\frac{-E}{\tilde{E}_0} = \frac{1}{(k+l+1)^2}$$

$$\boxed{E = \frac{-1}{(k+l+1)^2} \times \tilde{E}_0 = \frac{-1}{(k+l+1)^2} \times \frac{1}{2} Z^2 \alpha^2 \mu c^2}$$

The principal quantum number n is defined

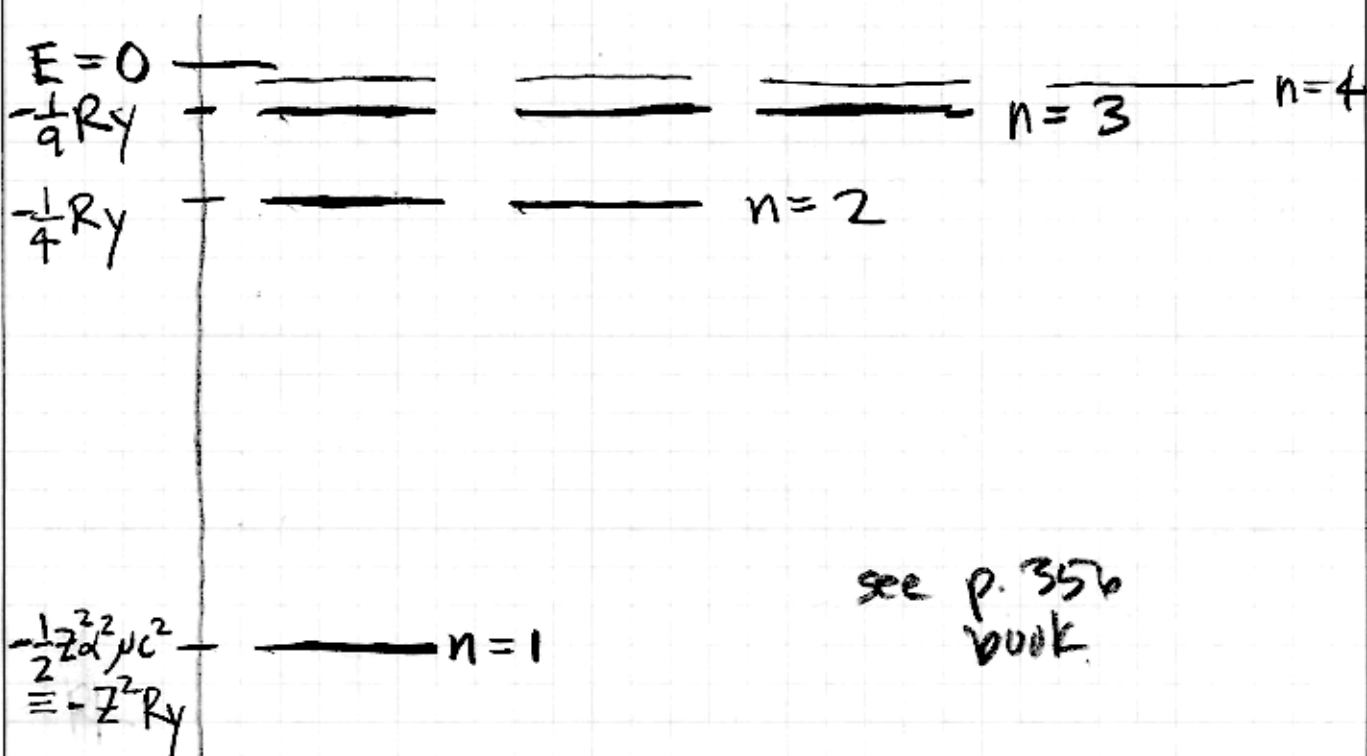
$$n = k + l + 1 \quad \text{minimum } n \text{ is then } \underline{0}$$

no maximum

But there is degeneracy; a variety of l can make the same n .

$$l = 0, 1, 2, \dots, n-1$$

The spectrum is often represented the following way:



$$Ry = \frac{1}{2} \alpha^2 \mu c^2 \quad \begin{array}{cccc} l=0 & l=1 & l=2 & l=3 \\ \text{"s"} & \text{"p"} & \text{"d"} & \text{"f"} \end{array}$$

note: an ∞ # of bound states
(remember?) since $|V(r)| \rightarrow 0$
slower than $1/r^2$

Wave Functionslots of backtracking
(p. 356 book)

$$R_{\ell\ell}(\tilde{r}) = \frac{U_{\ell\ell}(\tilde{r})}{\tilde{r}} = \tilde{r}^{\ell} e^{-\alpha \tilde{r}} W_{\ell\ell}(\tilde{r})$$

$$\alpha = \sqrt{-\tilde{E}} = \frac{1}{n} = \sqrt{-\frac{E}{E_0}} \quad \tilde{E}_0 = \frac{1}{2} Z^2 \mu c^2$$

$$W_{\ell\ell}(\tilde{r}) = \sum_{k=0}^{n-\ell-1} C_k \tilde{r}^k \propto \text{"Laguerre Polynomial"}$$

$$\propto L_{n-\ell-1}^{2\ell+1}(2\tilde{r})$$

remember: $\tilde{r} = \frac{r}{a} = \frac{Z}{a_0} r \quad a_0 = \frac{\hbar^2}{\mu e^2}$

$$L_p^0 = e^x \left(\frac{d^p}{dx^p} \right) e^{-x} x^p$$

$$p=0 = 1$$

$$= 1 \quad e^x [-e^{-x} + e^{-x}]$$

$$1-x$$

$$= 2 \quad e^x [e^{-x} - 2e^{-x} + 2e^{-x}]$$

$$2-2x+x^2$$

example: $k=1=2\ell+1, \ell=0$
 $n=1, p=0$

$$L_0^1 = (-1)^1 \frac{d}{dx} L_1^0 = -\frac{d}{dx} (1-x) = 1$$

$$k=2=2\ell+1, \ell=1$$

$$n=2, p=2-1-1=0$$

$$L_0^2 = (-1)^2 \frac{d^2}{dx^2} L_2^0 = \frac{d^2}{dx^2} (2-2x+x^2) = 2$$

... etc

limiting behavior: $\tilde{r} \rightarrow \infty$

$$R_{n\ell}(\tilde{r}) \sim \tilde{r}^{\ell} \times \tilde{r}^{n-\ell-1} e^{-\frac{\tilde{r}}{n}} \sim \tilde{r}^{n-1} e^{-\frac{\tilde{r}}{n}}$$

Normalizations - just look up.